

# Alternative ways to simulate a detector

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The 4th Electron Ion Collider Workshop

May 21, 2008

# Alternative to what? Monte Carlo and the detector design problem

Let us assume we are charged with the design of a large, multi-component, collider-scale detector.

- Can I do the physics? Overlapping issues:
  - 1 Are my statistics sufficient?
  - 2 Is the data rate tractable?
  - 3 Is my resolution adequate?
  - 4 Is my acceptance/efficiency large enough?
- Only complete answer for (3) & (4): complete Full-Blown GEANT-Style Monte Carlo (FBMC)

# What's involved in writing and using a FBMC?

## Tasks:

- detailed geometry
- hit generation
- digitization
- reconstruction
- event format and content issues
- computationally intensive
  - ▶ manage multiple jobs
  - ▶ multiple files

## Comments:

- a lot of work
- especially difficult for charged particle tracking
- desire: have this become part of permanent code base

Must all be done eventually: no substitute for full-blown Monte Carlo

# What if you want try out new ideas?

- detector technology choice
- optimizations
  - ▶ material budget
  - ▶ detector placement
  - ▶ resolution assumptions
- FBMC not practical or not possible (no infrastructure yet)
- Would be better if infrastructure to do the studies is cheap/disposable
- Sacrifice fidelity for speed (implementation and execution)

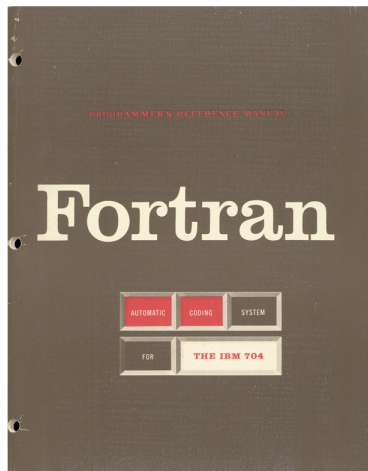
Look at two examples:

- 1 REZEST (resolution estimator)
- 2 MCFast (Monte Carlo, fast)

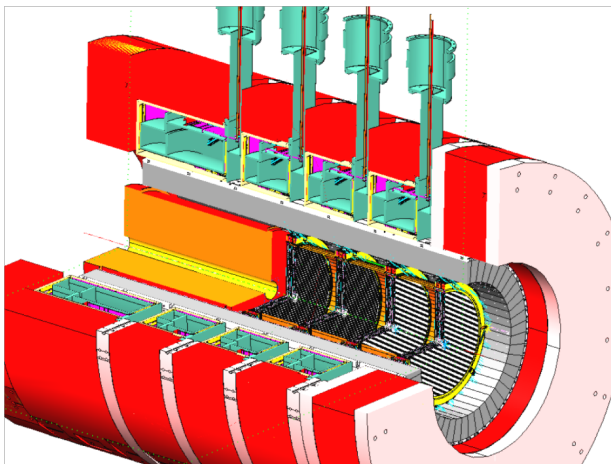
# What is REZEST?

The back-of-the-envelope, coded up

- A set of FORTRAN routines
- Charged track **re**zolution **est**imation in transverse momentum and direction for the GlueX geometry
- Use results as input to smearing routines
- Parameters can be varied to quickly obtain estimates for new configurations
- No Monte Carlo is used; results are returned immediately



# GlueX Detector Geometry



**CDC** : Central Drift Chamber, straw tubes, axial and stereo

**FDC** : Forward Drift Chamber, planar chambers,  $\perp$  to beamline

**Magnet** : Superconducting solenoid

# Approximations

The following assumptions are made:

- The magnetic field is uniform everywhere.
- Particles travel in straight lines, independent of momentum.
- All position measurements within a detector (FDC or CDC) are statistically independent of one another.
- All positions measurements within a detector are made at locations uniformly spaced along the trajectory.
- All positions measurements within a detector have the same resolution.

Relative variation of resolution when a particular parameters are varied should give a good feeling for the effect of parameter change.





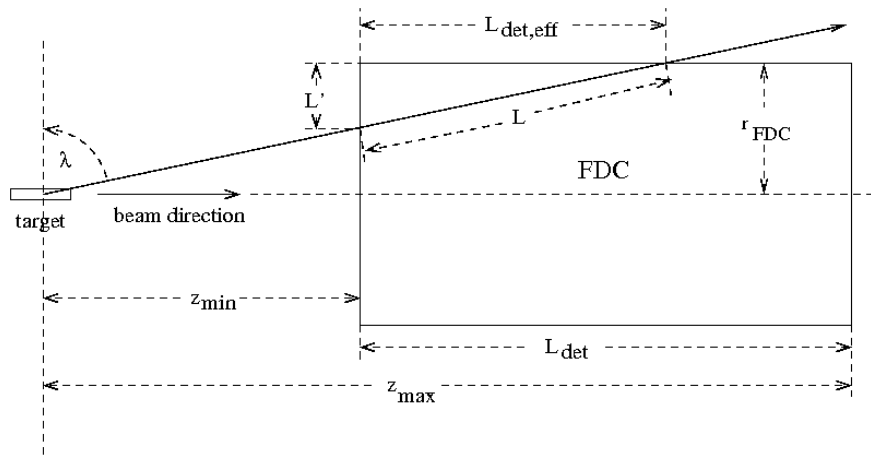
## CDC Geometry, 2

Other parameters, not shown are:

- $r_{\min,\text{stereo}}$  minimum radius of the CDC stereo layers
- $r_{\max,\text{stereo}}$  maximum radius of the CDC stereo layers
- $n_{\text{RL,CDC}}$  number of radiation lengths measured transverse to the tracking layers ( $n_{\text{r1}} = x/X_0$ )
- $n_{\text{RL,front}}$  number of radiation lengths in the material inside the CDC
- $n_{\text{RL,endplate}}$  number of radiation lengths in the downstream CDC endplate
- $n_{\text{m,CDC}}$  number of position measurements, total
- $n_{\text{m,stereo}}$  number of position measurements in stereo layers

Parameter	Value
$r_{\min}$	0.10960 m
$r_{\max}$	0.56534 m
$z_{\text{CDC}}$	1.02 m
$r_{\min,\text{stereo}}$	0.16304 m
$r_{\max,\text{stereo}}$	0.39473 m
$n_{\text{RL,CDC}}$	0.03437
$n_{\text{RL,front}}$	0.01437
$n_{\text{RL,endplate}}$	0.02810
$n_{\text{m,CDC}}$	25
$n_{\text{m,stereo}}$	8

# Geometry of the FDC



## Geometry of the FDC (2)

Other parameters are:

- $n_{\text{RL,FDC}}$  number of radiation lengths measured transverse to the tracking layers ( $n_{\text{rl}} = x/X_0$ )
- $n_{\text{m,FDC}}$  number of one-dimensional position measurements for a track which passes through all layers of the FDC

Parameter	Value
$z_{\text{min}}$	1.25 m
$z_{\text{max}}$	2.92 m
$r_{\text{FDC}}$	0.56534 m
$n_{\text{RL,FDC}}$	0.028258
$n_{\text{m,FDC}}$	24

## Transverse Momentum Resolution

The formulae used to estimate transverse are taken from the Particle Data Group's Review of Particle Physics. For a particle with charge  $q$  of momentum  $p$  in a uniform magnetic field  $B$  with a pitch angle  $\lambda$

$$p_t \equiv p \cos \lambda = (0.3)qBR \quad (1)$$

where  $R$  is the radius of curvature in the projection of the trajectory onto the bend plane,  $p$  is in GeV/c,  $B$  is in Tesla, and  $R$  is in meters. The curvature  $k = \frac{1}{R}$ . The variance of  $k$  has two contributions,

$$(\delta k)^2 = (\delta k_{\text{res}})^2 + (\delta k_{\text{ms}})^2. \quad (2)$$

$$\delta k_{\text{res}} = \frac{\epsilon}{L'^2} \sqrt{\frac{720}{N+4}} \quad (3)$$

where  $\epsilon$  is the position resolution in meters,  $L'$  is the projected length of the track onto the bending plane in meters and  $N$  is the number of measurements.

$$\delta k_{\text{ms}} = \frac{(0.016 \text{ GeV}/c)z}{Lp\beta \cos^2 \lambda} \sqrt{n_{\text{RL}}} \quad (4)$$

where  $n_{\text{RL}}$  is the number of radiation lengths in the detector and  $L$  is the total track length in the detector. For the momentum estimate, the amount of material in front of the detector is ignored.

## Error on Slope and $y$ -intercept of a Straight-Line Fit

To estimate the error due to position resolution on the direction of a fitted track, we use the error on the slope of a straight line fitted to the same number of measurements.

$$\chi^2 = \sum \left[ \frac{1}{\sigma_i^2} (y_i - a - bx_i)^2 \right] \quad (5)$$

The variance of  $b$  is

$$\sigma_b^2 \approx \frac{n\sigma^2}{\Delta'} \quad \text{where} \quad \Delta' = n \sum x_i^2 - \left( \sum x_i \right)^2 \quad (6)$$

For  $n$  equally spaced measurements spanning the interval  $[0, L]$ ,

$$x_i = \frac{L(i-1)}{n-1} \quad (7)$$

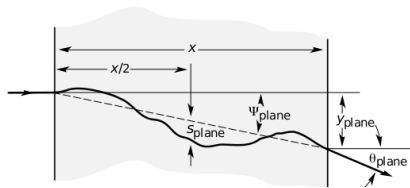
we get

$$\sigma_b^2 = \frac{12\sigma^2(n-1)}{L^2 n(n+1)} \quad (8)$$

We need to translate an error in slope to an error in angle.  $\theta = \tan^{-1} b$  so

$$\delta\theta = \left| \frac{d\theta}{db} \right| \delta b = \frac{\delta b}{\sec^2 \theta}. \quad (9)$$

# Angular Error Due to Multiple Coulomb Scattering



The central angular distribution is approximately Gaussian with a width given by

$$\theta_0 = \frac{(13.6 \text{ MeV})}{\beta c p} z \sqrt{x/X_0} [1 + 0.038 \ln(x/X_0)]. \quad (10)$$

The angle  $\Psi_{\text{plane}}$  is used as an approximation to the contribution of multiple scattering to both the azimuthal and polar angles.

$$\Psi_{\text{plane}}^{\text{rms}} = \frac{1}{\sqrt{3}} \theta_0. \quad (11)$$

The material in front of a particular detector (“fronting material”) is included as an addition to the number of radiation lengths in the detector itself.

# Contribution to Azimuthal Angle Resolution from Curvature Resolution

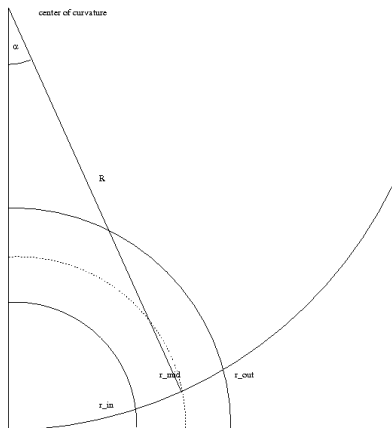
To infer the azimuthal angle  $\phi$  at the vertex, track must be swum backward through angle  $\alpha$ .

$$\sin \frac{\alpha}{2} = \frac{r_{\text{mid}}}{2R} = \frac{r_{\text{mid}}k}{2} \quad (12)$$

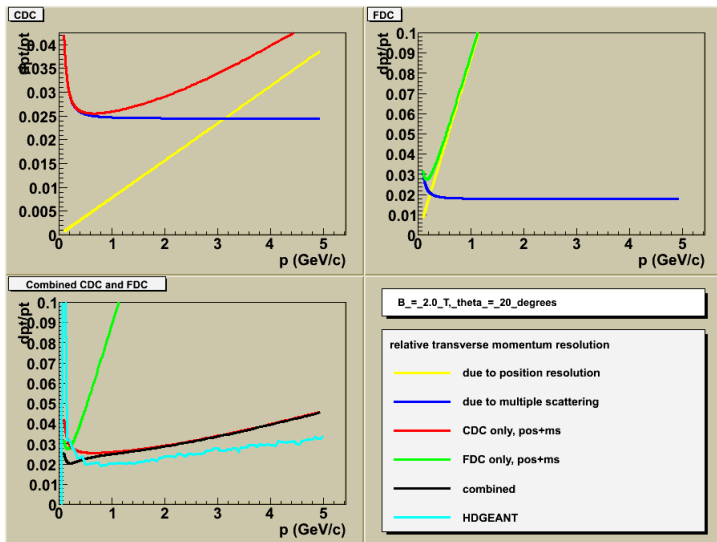
so

$$\delta\alpha = r_{\text{mid}}\delta k \sec \frac{\alpha}{2} \quad (13)$$

where  $\alpha = 2 \sin^{-1}(r_{\text{mid}}/2R)$ . As an approximation, we take  $r_{\text{mid}} = (r_{\text{in}} + r_{\text{out}})/2$ .

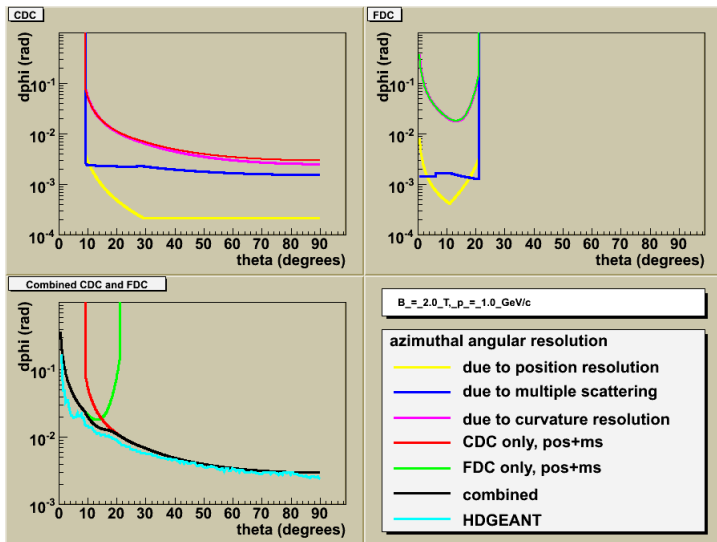


# Resolution in relative transverse momentum as a function of total momentum at $20^\circ$ for $B = 2.0$ T.





# Resolution in azimuthal angle as a function of polar angle at $p = 1.0 \text{ GeV}/c$ for $B = 2.0 \text{ T}$ .



# Conclusions on REZEST

- ① The plots show reasonable agreement with the HDGEANT (the GlueX FBMC) results. Agreement is generally at the 20% level, in some places better, in others as poor as a factor of 2.
- ② Rather detailed features of resolution variation are exhibited faithfully.
- ③ One area where the simple model can break down is in the straight-line approximation for the trajectories for particles with very low transverse momentum.
- ④ Some acceptance information available by excluding poor resolution regions.
- ⑤ The most profitable use in predicting relative changes in resolution as detector parameters are changed.

# What is MCFast?

physics features:

- charged particle tracking
  - ▶ position resolution
  - ▶ multiple scattering
  - ▶ energy loss
- calorimetry
  - ▶ em
  - ▶ hadronic
  - ▶ parametrized showers

technical features:

- built-in interface to common event generators (pythia, qq)
- creates event stream: true Monte Carlo
- detector geometry specified in an ascii file
- hooks for user intervention and event examination

# History of MCFast

- developed for B-TeV design studies at FNAL
- 1994, v1.4, wrapper around SLAC TRACKERR program
- 1995, v2.1, complete rewrite
- ca. 2001, v5.2, recommended version
- significant manpower investment...
- ...but “not supported” anymore
- 2001 GlueX/Hall-D effort, important for initial design studies

# MCFast Geometry Specification



**R-plane model** right circular cylinders or polygonal shells, centered on z-axis, with defined radius, z center and z length

**Z-plane model** planes perpendicular to z axis, rectangular or circular outer boundary, may have beam hole

**Conical model** cones centered on z axis

**Special cases** planes perpendicular to x or y for magnet yokes, calorimeters (not for tracking volumes)

All volumes have material composition specified.

# MCFast charged particle tracing and “fitting”

## definitions

**tracing**: stepping particle through material and magnetic fields

**“fitting”**: obtaining track parameters and covariance matrix

# Tracing

Events recorded:

- 1 radial plane encountered
- 2 z plane encountered
- 3 conical surface encountered
- 4 production point
- 5 decay in flight
- 6 pair production
- 7 absorption
- 8 E & M shower starts
- 9 hadron shower starts
- 10 dummy point (for display purposes)

Treatment of magnetic fields:

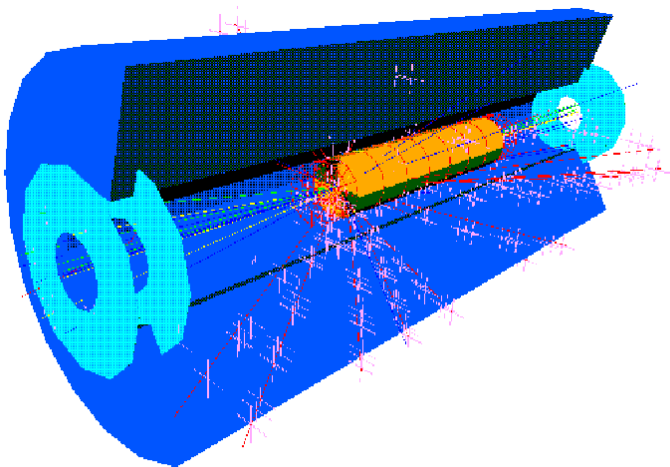
- define regions of constant magnetic field
- non-uniform fields must be approximated by discrete regions of constant field

Energy loss and multiple scattering can be enabled separately (or together).

# “Fitting”

- no pattern recognition
- use hits from the tracing
- minimum hit requirement simulates detector acceptance
- two approaches:
  - 1 pseudo fitting
    - ★ with or without energy loss, MCS
    - ★ calculates covariance matrix (CM) among track parameters
    - ★ uses CM to smear track parameters
    - ★ no  $\chi^2$  computed
  - 2 Kalman filter
    - ★ full Kalman filter algorithm
    - ★ energy loss and/or multiple scattering accounted for (if turned on in tracing)
    - ★ CM calculated
    - ★ smeared track can be generated
    - ★  $\chi^2$  generated





## MCFast model of D0 tracking upgrade

# Why is MCFast different than FBMC?

- ① simplified geometry specification
- ② simplified magnetic field map
- ③ no pattern recognition
- ④ no event serialization
- ⑤ pseudo fit much faster than a full fit

# Conclusions on MCFast

- ① Within the limits of the simple geometry specification, it gives a highly realistic rendering of detector response.
- ② It represents a not insignificant implementation effort.
- ③ Lack of current support is a concern.

# General conclusions

- 1 In the early stages of the design cycle, a quick parameter-based estimation of detector performance is important.
- 2 Two examples have been discussed: REZEST and MCFast.
- 3 Usefulness of tools like this will likely extend beyond design phase. (What if's always seem to come up.)
- 4 It is worthwhile spending some time developing these types of facilities.

Acknowledgements: Paul Eugenio (Florida State), Lynn Garren and Patricia McBride (Fermilab)

References:

- 1 REZEST:  
<http://argus.phys.uregina.ca/cgi-bin/public/DocDB/ShowDocument?docid=1015>
- 2 MCFast: <http://cepa.fnal.gov/psm/simulation/mcfast/>