

# The Lead/Scintillating Fibre Barrel Calorimeter (BCAL)

Blake Leverington

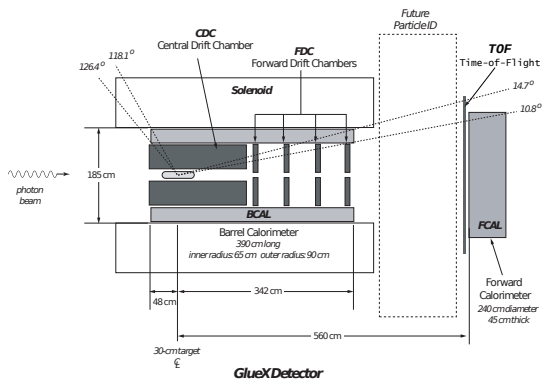
Department of Physics  
University of Regina

May Colloboration meeting  
JLAB  
May 13, 2010

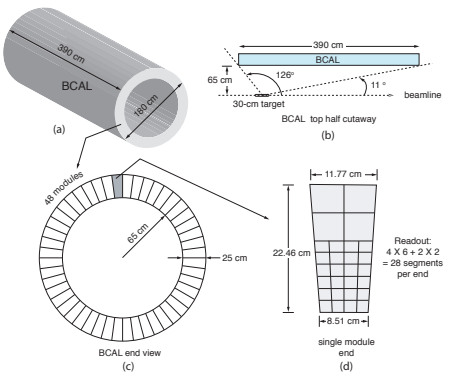
# Outline

- 1 The Barrel Calorimeter in GlueX
- 2 Module Details
- 3 Fibre Details
- 4 Electromagnetic Showers
- 5 Resolutions
- 6 The BCAL in the Monte Carlo

# The BCAL in GLUEX

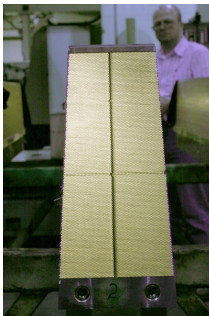


**Figure:** Schematic of the GlueX Detector. The detector has cylindrical symmetry about the beam direction.

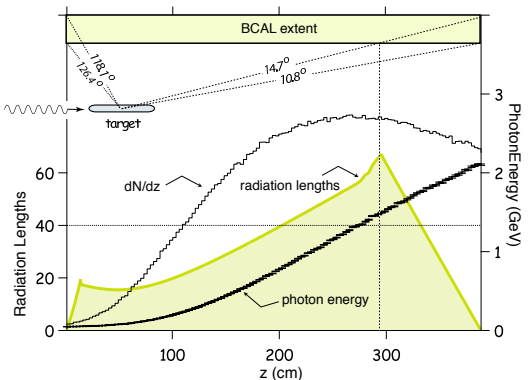


**Figure:** The GlueX BCAL. (a) BCAL schematic; (b) a BCAL module side view; (c) end view of the BCAL showing all 48 modules and (d) an end view of a module showing readout segmentation.

The calorimeters play multiple roles: Primarily, they are to detect the photons from  $\eta$  and  $\pi^0$  decays; The positions and energies of the photons must be determined with sufficient accuracy and resolution resolution to allow for a complete kinematic reconstruction of the event. Secondly, some dE/dx and PID information.

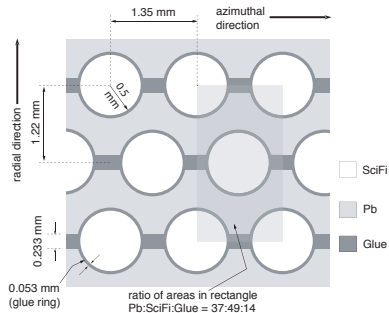


**Figure:** A photograph of a final production module for the BCAL.

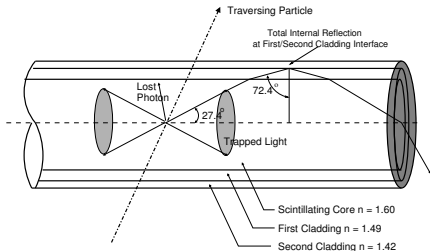


**Figure:** The distribution of photons, their energy and integrated path length through the Pb/SciFi matrix as a function of position along the length of the BCAL for one of the GlueX signature reactions,  $\gamma p \rightarrow \eta \pi^0 p \rightarrow 4\gamma p$ , is shown. The target position and angular range subtended by the BCAL are also presented.

# Module Details



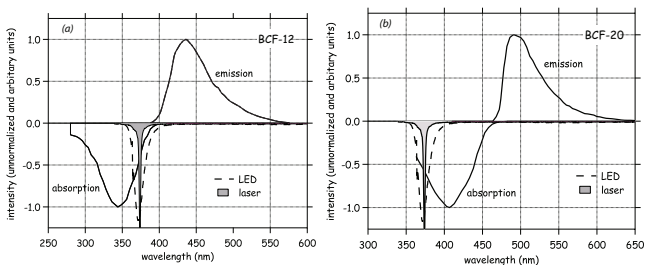
**Figure:** The BCAL fibre matrix showing the placement of 1 mm diameter fibres in the azimuthal and radial directions. Particle tracks would appear to enter the matrix from the bottom.



**Figure:** A simple schematic of a standard double-clad fibre showing the trapping of optical light for meridional rays.

- 0.96 mm core of doped polystyrene and two layers of polymethylmethacrylate cladding: acrylic (3%) and fluor-acrylic material (1%).
- the passage of ionizing electrons through the core produce 8000 optical photons/MeV (disputed)



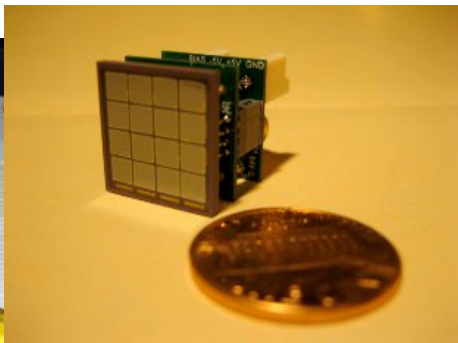
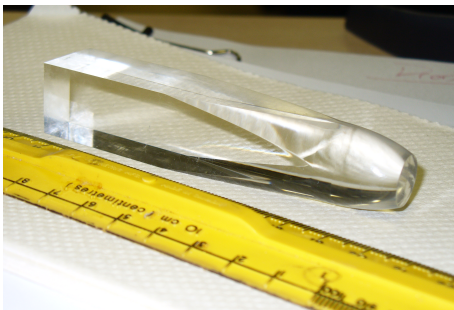


**Figure:** Emission and absorption spectra from the secondary dye of (a) blue BCF-12 and (b) green BCF-20 fibres.

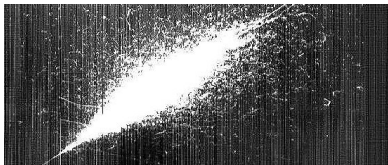
- Kuraray SCSF-78MJ fibres are used for the production modules
- each 1mm fibre will trap  $\sim 5.3\%$  (min.) (total internal reflection) of the produced light in the fibre
- the fibres have an average bulk attenuation length of  $\sim 350$  cm.
- the readout detectors must be sensitive to the optical spectrum of the fibres

# Readout

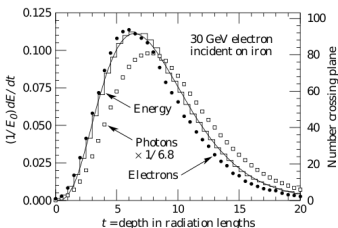
- Attached to the end of each module is an array of acrylic light-guides to concentrate the light to a smaller area which is read-out a Silicon Photomultiplier array.
- The large amount of noise from the SiPM sets the threshold for the energy in each BCAL cell.



# Electromagnetic Showers

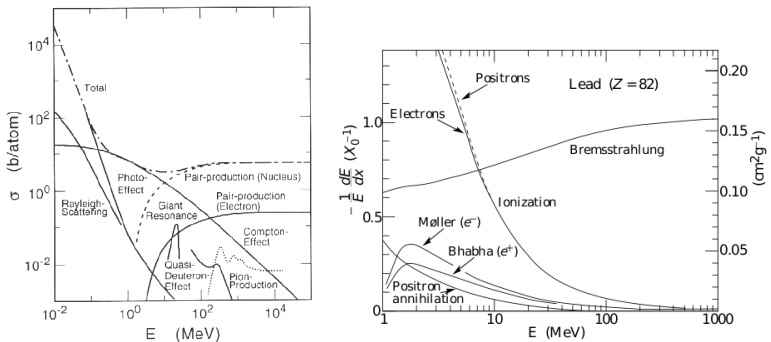


Reminder: radiation length is both (a) the mean distance over which a high-energy electron loses all but  $1/e$  of its energy by bremsstrahlung, and (b)  $7/9$  of the mean free path for pair production by a high-energy photon



**Figure:** The fractional energy deposition per radiation length.

Total thickness of the BCAL perpendicular to the beam :  $15.5X_0$  where  $X_0 = 7.06 \text{ g/cm}^2$  or  $1.45 \text{ cm}$ .



**Figure:** (left) Photon total cross sections as a function of energy in lead. (right) Fractional energy loss per radiation length in lead as a function of electron or positron energy. Electron (positron) scattering is considered as ionization when the energy loss per collision is below 0.255 MeV.

- Note: A small fraction (0.2%) of the bremsstrahlung energy will produce hadrons.

# Energy Resolution

- The BCAL is a sampling calorimeter: an active medium (Sci.Fi.) which generates signal and a passive medium as an absorber (lead).
- the stochastic term  $a$  for a sampling calorimeter is expected to be proportional to  $\sqrt{t/f}$ , where  $t$  ( $\sim 0.5\text{mm}$ ) is the absorber thickness and  $f$  is the sampling fraction ( $f = 12.5\%$ ).
- Intrinsic shower fluctuations result in a stochastic term in the energy resolution: due to sampling fluctuations ( $\sim 4\%$ ), photoelectron statistics ( $\sim 2\%$ ), and dead material in front of the calorimeter.
- Detector non-uniformity and calibration uncertainty result in a systematic, or constant, term.

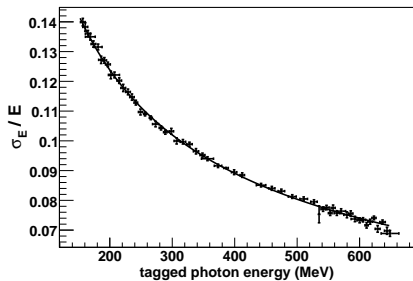


Figure: Energy resolution vs.  $E_{\text{BEAM}}$  for photons for  $\theta = 90^\circ$  and  $z = 0$  cm.

From JLAB2006 photon beam test in Hall-B:

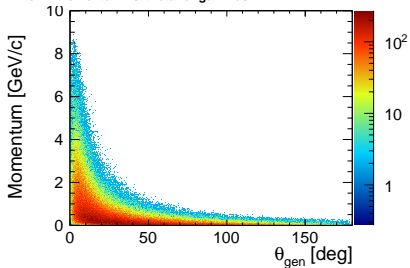
$$\frac{\sigma_E}{E} = \frac{5.4\%}{\sqrt{E(\text{GeV})}} \oplus 2.3\%, \quad (1)$$

# Timing Resolution

- the double ended readout allows for the position of the shower to be determined
- a time difference resolution from the beam test of  $\sigma_{\Delta T/2} = 70$  ps/ $\sqrt{E(\text{GeV})}$  ps.
- this results in a position resolution of  $\sigma_z = 1.1$  cm for a 1 GeV photon.
- energy and position resolution agree with KLOE's beam test results.

# Monte Carlo: What do we expect?

Thrown momentum vs. theta for gammas



Pythia Photons

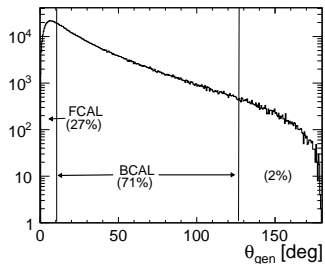
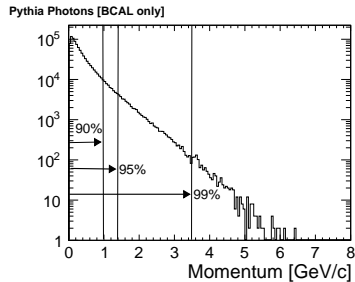


Figure: (left) Generated momentum vs.  $\theta$  distribution for photons from PYTHIA decays. (right) The projection of (a) onto the  $\theta$  axis.



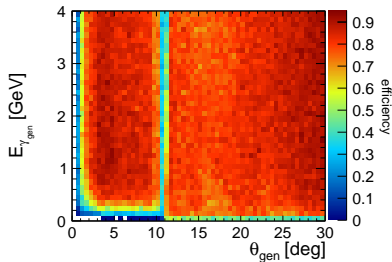
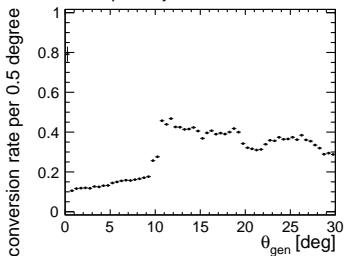


**Figure:** The momentum distribution for photons which would strike the BCAL only.

# The BCAL in the Monte Carlo

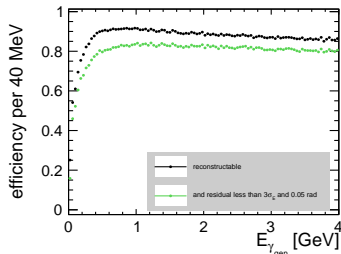
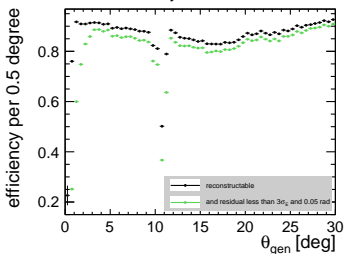
- Energy and timing resolutions from the beam test have been encoded in HDGEANT along with sampling fluctuations (from other Monte Carlo) and expected SiPM detector noise.
- Cells with hits close by that exceed a threshold are clustered together to create a “DPhoton” object.
- How does the other detector material affect the BCAL?
- Generate photons with a range of energy (0 to 4 GeV) and illuminate the entire BCAL!

Photon conversion probability



**Figure:** (left) Single photon conversion probability as a function of incident polar angle from  $0^\circ$  to  $30^\circ$ . (right) Single photon reconstruction efficiency as a function of energy and polar angle.

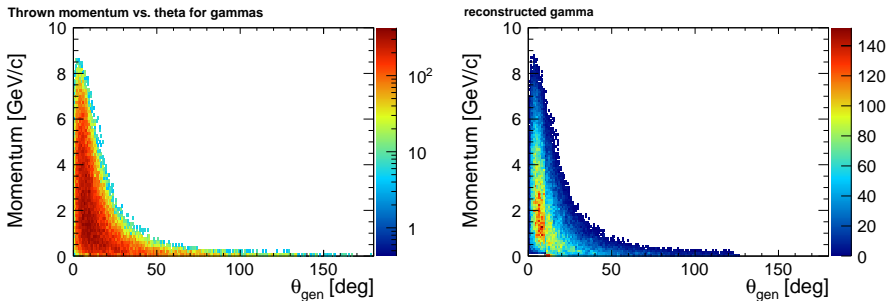
Photon reconstruction efficiency



**Figure:** (left) The polar angle projection from previous slide. The efficiency clearly degrades in the gap region. (right) The energy axis projection from previous slide.

let's try an  $\eta\pi^0 \rightarrow 4\gamma$ !

What if we use everything we've learned and try to reconstruct  $\eta\pi^0 \rightarrow 4\gamma$ .



**Figure:** The momentum versus polar angle distribution for photons from  $\eta\pi^0 \rightarrow 4\gamma$  decays. (left) Thrown events and (right) reconstructed events shown with a linear scale to show the gap in efficiency at  $\theta_{\text{gen}} = 10^\circ$  more clearly.

# Not everything is perfect...

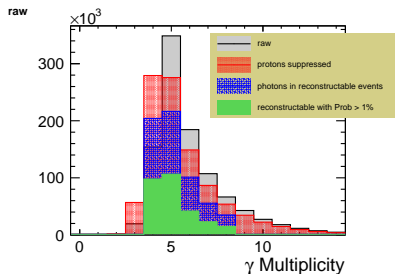


Figure: Photon multiplicity for reconstructed  $\eta\pi^0 \rightarrow 4\gamma$  signal events.

- Extra clusters are formed in many cases (split-offs, conversions in other detectors) Messy!
- Events with charged particles are even messier, as they will also form clusters.

# Summary

- the final BCAL modules are currently being constructed in Regina
- the properties of the BCAL are well understood and are included in MC
- the BCAL seems to function well in MC to reconstruct  $\eta\pi^0$  decays
- however, some improvements in reconstruction are still needed to identify photon showers properly with the BCAL