A great deal of early information can be found in the first edition of the Particle Data Book [1] published in 1964, and the subsequent version [2] published in 1965.

- The $S(975)$ or the S^* is now known as the $f_0(980)$. The S^* was reported in 1966 [3]. The state was first reported in the reaction $\pi^+ p \to \Delta^{++}\pi^+\pi^-$ and $\pi^+ p \to \Delta^{++}K^+K^-$ in 1973 [4].

- The $\delta(980)$ is now known as the $a_0(980)$. The state was first reported in $\bar{p}p \to K_S K^{\pm}\pi^{\mp}$ at rest, where the analysis focused on events without $K^*(890)$ in 1965 [5].

- The $H(1190)$ is now known as the $h_1(1170)$. The state was first observed in $\pi^- p \to \pi^+\pi^-\pi^0 n$ in 1981 [6].

- The $B(1235)$ is now known as the $b_1(1235)$. The state was first observed in an analysis of reactions of the type $\pi^+ p \to pX$ in 1963. [7]

- The $f(1270)$ is now known as the $f_2(1270)$. The state was first observed in $\pi^- p \to n\pi^+\pi^-$ in 1963 [8].

- The $A(1270)$ or the A_1 is now known as the $a_1(1260)$. An enhancement in the region of the a_1 mass was first observed in π^- interactions on carbon. The authors separated their data into classes of events based on the number of pions. The so-called A region contained $\pi^+\pi^-\pi^-$ events and a broad enhancement was seen [9]. The A name stuck for mesons in this mass region. In subsequent work [10, 11], two states were resolved: the A_1 and A_2.

- The $D(1285)$ is now known as the $f_1(1285)$. The state was first reported in $\pi^- p$ interactions in three or more-body final states. It was first reported in 1967 [12]

- The $\epsilon(1300)$ is know known as the $f_0(1370)$. The state was first reported in $\bar{p}n \to \pi^+\pi^+\pi^-\pi^-\pi^-$ annihilation at rest as an enhancement in $\rho\rho$. It was first reported in 1966 [13].

- The $A_2(1320)$ is now known as the $a_2(1320)$. This appears to be firest reported in πp interactions, but I have been unable to track down the earliest reference. It was orignially discussed as part of the A which ultimately included both the A_1 and A_2 [9, 10, 11].

- The $E(1420)$ is now known as the $f_1(1420)$. The $E(1420)$ and the $\iota(1440)$ have a bit of a confused history. A state called the E with $J^{PC} = 0^{-+}$ was first reported in $\bar{p}p \to K\bar{K}3\pi$ in the $K\bar{K}\pi$ invarant mass [14]. The E was first reproted oin $J^{PC} = 0^{-+}$ was reported in radiateive J/ψ decays in $K\bar{K}\pi$ [15]. A $J^{PC} = 1^{++}$ state was later observed in πp inetarcations [16] and by 1984, the E was assigned $J^{PC} = 1^{++}$ quantum numbers and the ι was assigned $J^{PC} = 0^{-+}$ quantum numbers. Hence, older literature will likely be very confused on these two states. The pseudoscalar states are prodcued strongly in both radaitive J/ψ decays and proton antiproton annilation at rest. The axial state is produced in central production, photoproduction and πp interactions.

- The $\iota(1440)$ is now two states, the $\eta(1405)$ and the $\eta(1475)$.

- The $f'(1525)$ is now known as the $f_2'(1525)$. This state was first observed in $K^- p$ unteractions in ${}_S K_S$ recoiling against a hyperon. The original observation is from 1965 [17] and a spin parity analysis from 1966 [18].

- The $\omega(1670)$ is now known as the $\omega_3(1670)$.
- The $A(1680)$ or A_3 is now known as the $\pi_2(1670)$.
- The $\theta(1690)$ is now known as the $f_0(1700)$. There was early uncertainty whether it was $J^{PC} = 0^{++}$ or 2^{++}. It is now confirmed as 0^{++}.
- The $\phi(1850)$ is now known as the $\phi_3(1850)$.
- The $h(2030)$ is now known as the $f_4(2050)$.
- The $Q(1280)$ or Q_1 is now known as the $K_1(1270)$.
- The $\kappa(1350)$ is now known as the $K_0^*(1430)$.
- The $Q(1400)$ or Q_2 is now known as the $K_1(1400)$.
- The $K^*(1430)$ is now known as the $K_2^*(1430)$
- The $L(1770)$ is now known as the $K_2(1770)$.
- The $K^*(1780)$ or K^* is now known as the $K_3^*(1780)$.
- The $K^*(2060)$ is now known as the $K_4^*(2045)$.

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