

Electromagnetic calorimeters based on scintillating lead tungstate crystals for experiments at Jefferson Lab [☆]

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Abstract

A new electromagnetic calorimeter consisting of 140 lead tungstate (PbWO₄) scintillating crystals was constructed for the PrimEx- η experiment at Jefferson lab. The calorimeter was integrated into the data acquisition and trigger systems of the GlueX detector and used in the experiment to reconstruct Compton scattering events. The experiment started collecting data in the spring of 2019 and acquired about 30% of the required statistics. The calorimeter is a prototype for two PbWO₄-based detectors: the Neutral Particle Spectrometer (NPS) and the lead tungstate insert of the forward calorimeter (FCAL) of the GlueX detector. The article presents the design and performance of the Compton calorimeter and gives a brief overview of the FCAL and NPS projects.

Keywords: Electromagnetic calorimeter, Lead tungstate scintillator

1. Introduction

Electromagnetic calorimeters based on PbWO₄ scintillating crystals have a widespread application in experiments at different accelerator facilities such as CERN, FNAL, GSI, and Jefferson Lab (JLab). The small radiation length ($L_R = 0.89$ cm) and Molière radius ($R_M = 2.19$ cm) of PbWO₄ allows to build high-granularity detectors with a good spatial separation and energy resolution of reconstructed electromagnetic showers, which makes these crystals the material of choice in many of these applications.

Two electromagnetic calorimeters are currently under construction in experimental Hall D and Hall C at Jefferson Lab, both using rectangular 2.05 cm \times 2.05 cm \times 20 cm PbWO₄ scintillating modules. The inner part of the forward lead glass calorimeter of the GlueX detector [1] in Hall D will be upgraded with these high-granularity, high-resolution crystals. This upgrade is required by the JLab Eta Factory (JEF) experiment to perform precision measurements of various $\eta^{(\prime)}$ decays with emphasis on rare neutral modes [2]. The Neutral Particle Spectrometer [3] in experimental Hall C consists of a PbWO₄ electromagnetic calorimeter preceded by a sweeping magnet. The

NPS is required by Hall C's precision cross section measurement program with neutral final states [4–9]. Such precision measurements of small cross sections play a central role in studies of transverse spatial and momentum hadron structure. The NPS detector consists of 1080 PbWO₄ crystals arranged in a 30 \times 36 array. Lead tungstate crystals for both detectors were procured from two vendors: Shanghai Institute of Ceramics (SICCAS) in China and CRYTUR in the Czech Republic. The quality of recently produced PbWO₄ scintillators has been studied in detail by the NPS and EIC eRD1 collaborations and is described in Ref. [10]. PbWO₄ crystals are also being considered for an electromagnetic calorimeter of the future Electron-Ion Collider [11].

In this article we describe the design and construction of a calorimeter prototype composed of 140 SICCAS crystals, which served as the Compton Calorimeter (CCAL) in the PrimEx- η experiment [12] with the GlueX detector in the spring of 2019. The CCAL was subsequently used during a few short GlueX physics runs at high luminosity in order to study rates and operating conditions expected for the FCAL lead-tungstate insert. Experience gained during fabrication and operation of the CCAL was critical for finalizing the design of the FCAL insert and also helped further optimize the NPS calorimeter.

This article is organized as follows: we will present the PrimEx- η experiment and performance of the CCAL in Section 2 and Section 3, and will briefly describe the FCAL and NPS projects in Sections 4 and 5.

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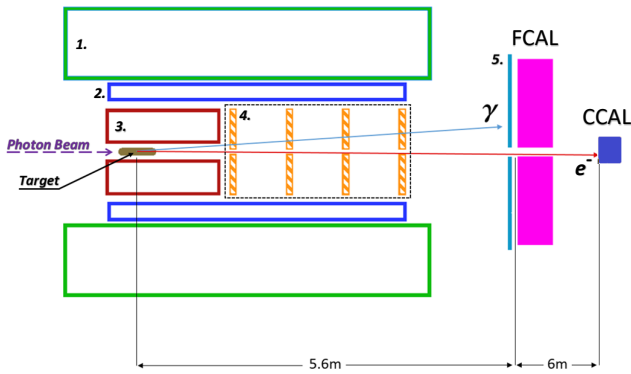


Figure 1: Schematic layout of the GlueX detector (not to scale). Numbers represent the following detector components: solenoid magnet (1), barrel calorimeter (2), central drift chamber (3), forward drift chambers (4), time-of-flight wall (5).

2. PrimEx- η experiment with the GlueX detector

The GlueX detector [1] was designed to perform experiments using a photon beam. Beam photons are produced via the bremsstrahlung process by electrons, provided by the JLab electron accelerator facility, incident on a thin radiator. The energy of a beam photon (E_γ) is determined by detecting a scattered electron after radiating the photon as follows: $E_\gamma = E_e - E'_e$, where E_e is the primary electron beam energy and E'_e is the energy of the bremsstrahlung electron. The bremsstrahlung electron is deflected in a 6 m long dipole magnet operated at a field of ~ 1.5 T and registered in the so-called tagging scintillator counters. Each counter corresponds to the specific energy of the reconstructed lepton. The tagging detectors span the beam photon energy range between 25% and 98% of the electron beam energy and covered the range between 2.8 GeV and 11.0 GeV during the PrimEx- η experiment¹. The typical energy resolution of the beam photon is about 0.1%. The photon beam propagates toward the GlueX target. A schematic view of the GlueX detector is illustrated in Fig. 1².

The physics goal of the PrimEx- η experiment is to perform a precision measurement of the $\eta \rightarrow \gamma\gamma$ decay width. The measurement will provide an important test of quantum chromodynamics symmetries and is essential for the determination of fundamental properties such as the ratios of the light quark masses and the η - η' mixing angle. The decay width will be extracted from the measurement of the photoproduction cross section of η mesons in the Coulomb field of a nucleus, which is known as the Primakoff effect. The η mesons will be reconstructed by

¹The electron beam energy during most production PrimEx- η runs was 11.2 GeV.

²Not shown on this plot is the DIRC detector, which was installed after the PrimEx- η experiment and is used for the particle identification in the forward direction.

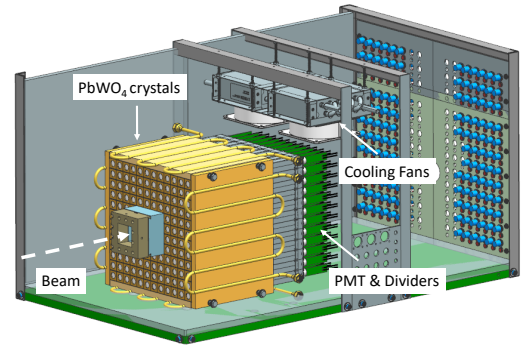


Figure 2: Schematic layout of the Compton calorimeter.

detecting two decay photons in the forward calorimeter of the GlueX detector.

The cross section will be normalized using the Compton scattering process, which will also be used to monitor the luminosity and control the detector stability during data taking. Electrons and photons originating from Compton events in the target are produced at small angles, typically outside the acceptance of the FCAL. In order to improve the reconstruction of particles in the forward direction, we built a small Compton calorimeter consisting of 140 lead tungstate scintillating crystals and positioned it about 6 m downstream from the FCAL as shown in Fig. 1. The CCAL covers the angular range between 0.19° and 0.47° .

The PrimEx- η experiment started collecting data in the spring of 2019 and has acquired 30% of the required statistics. During the experiment, the magnetic field of the solenoid magnet was switched off in order to allow reconstruction of Compton events. The photon flux was about $5 \cdot 10^6$ γ /sec (about five times lower than the nominal GlueX flux) in the beam energy range of interest between 9.5 GeV and 11.6 GeV.

3. Compton calorimeter of the PrimEx- η experiment

3.1. Calorimeter design

The calorimeter design is shown in Fig. 2. The CCAL comprises an array of 12×12 lead tungstate modules with a 2×2 hole in the middle for the passage of the photon beam. The modules are positioned inside a light tight box. A tungsten absorber is placed in front of the innermost layer closest to the beamline to provide protection from the high rate of particles predominantly originating from electromagnetic interactions.

The light yield from PbWO₄ crystals depends on temperature with a typical coefficient of $2\%/^\circ C$ at room temperature. Maintaining constant temperature is essential for the calorimeter operation. The calorimeter modules are surrounded by four copper plates with built-in pipes to circulate a cooling liquid and provide temperature stabilization. Foam insulation surrounds the detector box. The temperature was monitored and recorded during the experiment by five thermocouples attached to different points of the PbWO₄ module assembly. During the experi-

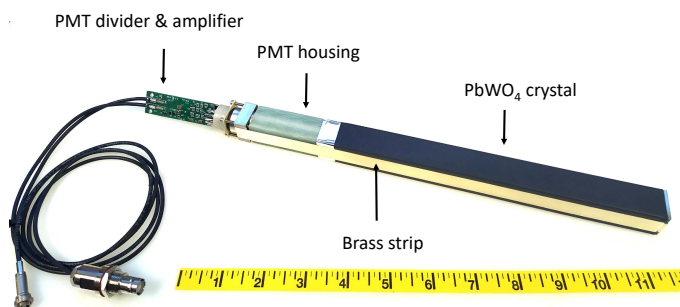


Figure 3: Calorimeter module showing main components: the PbWO₄ crystal, PMT housing, PMT divider, and signal and high-voltage cables.

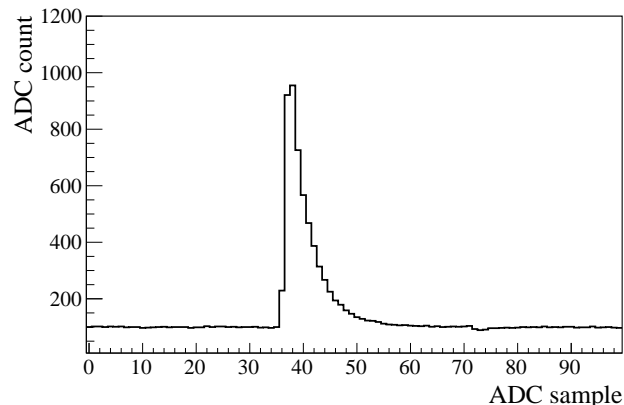


Figure 4: A typical flash ADC signal pulse obtained from a PbWO₄ module.

115 ment the temperature was maintained at $17^\circ \pm 0.2^\circ\text{C}$. The typ-151
 116 ical heat released by the photomultiplier tube (PMT) dividers
 117 of the whole detector was equivalent to about 30 Watts. In or-152
 118 der to prevent condensation, a nitrogen purge was applied. Two153
 119 fans with a water-based cooling system were installed on the154
 120 top of the crystal assembly to improve nitrogen circulation and155
 121 heat dissipation from the PMT dividers. The detector was po-156
 122 sitioned on a platform, which allowed to move it in the vertical157
 123 and horizontal directions, perpendicular to the beam. The plat-158
 124 form was remotely controlled and provided a position accuracy159
 125 of about $200\ \mu\text{m}$. During detector calibration each module was160
 126 moved into the beam.

127 3.2. Module design

128 The design of the PbWO₄ module is based on the HyCal167
 129 calorimeter, which was used in several experiments in Jeffer-168
 130 son Lab Hall B [13, 14]. An assembled calorimeter module169
 131 is presented in Fig. 3. Each lead tungstate crystal is wrapped170
 132 with a $60\ \mu\text{m}$ polymer Enhanced Specular Reflector film (ESR)171
 133 manufactured by 3MTM, which allows 98.5% reflectivity across172
 134 the visible spectrum. In order to improve optical isolation of173
 135 each module from its neighbors, each crystal is wrapped with a174
 136 layer of $25\ \mu\text{m}$ thick Tedlar. The PMT is located inside a G-10175
 137 fiberglass housing at the rear end of the crystal. Two flanges are176
 138 positioned at the crystal and housing ends and are connected to-177
 139 gether using $25\ \mu\text{m}$ brass straps, which are brazed to the sides178
 140 of the flanges. Four set screws are pressed to the PMT housing179
 141 flange to generate tension in the straps and hold the assembly180
 142 together. Light from the crystal is detected using a ten-stage181
 143 Hamamatsu PMT 4125, which is inserted into the housing and182
 144 is coupled to the crystal using optical grease (EJ-550). The183
 145 PMT diameter is 19 mm. The PMT is pushed towards the crys-184
 146 tal by using a G-10 retaining plate attached to the back of the185
 147 PMT and four tension screws applied to the PMT flange. The186
 148 PMT is instrumented with a high-voltage (HV) divider and am-187
 149 plifier positioned on the same printed circuit board attached to188
 150 the PMT socket.

3.3. Electronics

161 The PMT of each calorimeter module was equipped with an
 162 active base prototype [15], which was designed for the Neutral
 163 Particle Spectrometer in experimental Hall C. The base com-
 164 bines a voltage divider and an amplifier powered by the current
 165 flowing through the divider. The active base allows the oper-
 166 ation of the PMT at lower voltage and consequently at lower
 anode current, which improves the detector rate capability and
 prolongs the PMT's life. The original Hamamatsu divider for
 this type of PMT was modified by adding two bipolar transis-
 tors on the last two dynodes, which provides gain stabilization
 at high rate. The active base has a relatively large amplification
 of about a factor of 24 due to the large PMT count rate pre-
 dicted by Monte Carlo simulation of the NPS detector. Large
 amplification was not needed for the planned run conditions
 of the PrimEx- η experiment. However, we subsequently used
 CCAL in GlueX runs at significantly larger luminosity in order
 to study run conditions of the FCAL lead tungstate insert,
 where the amplifier will be required. This will be discussed in
 Section 4.0.3. During the PrimEx run, the CCAL PMTs were
 operated at about 680 V, which produced a divider current of
 $260\ \mu\text{A}$. The high voltage for each PMT was supplied by a
 24-channel CAEN A7236SN module positioned in a SY4527
 mainframe.

Amplified PMT signals were digitized using a twelve-bit 16-
 channel flash ADCs electronics module operated at a sampling
 rate of 250 MHz. The ADC was designed at Jefferson Lab [16]
 and is used for the readout of several sub-detectors of the GlueX
 detector. The Field-Programmable Gate Array (FPGA) chip
 inside the ADC module allows the implementation of various
 programmable data processing algorithms for the trigger and
 readout. An example of a flash ADC signal pulse obtained
 from a calorimeter module is shown in Fig. 4. In this exam-
 ple, the ADC is operated in the raw readout mode, where digi-
 tized amplitudes are read out for 100 samples, corresponding
 to the read out window size of 400 ns. During the PrimEx- η
 experiment, the ADC performed on-board integration of signal
 pulses, which amplitudes were above a threshold of 24 MeV.
 Amplitudes were summed in a time window of 64 ns and read

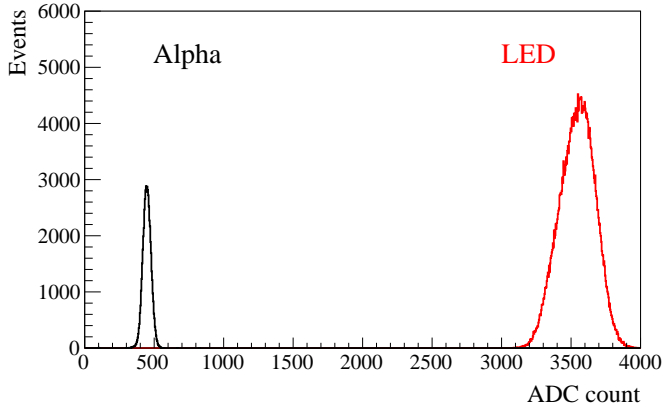


Figure 5: Flash ADC signal amplitudes induced by the LED and α -source in the reference PMT.

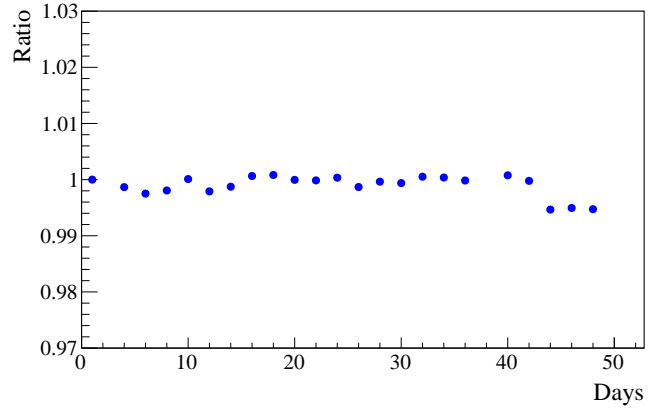


Figure 6: Ratio of signal ADC amplitudes from the LED pulser to the α -source measured by the reference PMT during different run periods of the 48-day long PrimEx- η experiment. The ratio is normalized to data in the beginning of the run.

190 out from the ADC module along with other parameters such
 191 as the pulse peak amplitude, pulse time, and data processing
 192 quality factors. This readout mode allowed to significantly re-
 193 duce the data size and ADC readout time, and therefore did not
 194 induce any dead time in the data acquisition.

195 CCAL flash ADCs are positioned in a VXS (ANSI/VITA
 196 41.0 standard) crate. VXS crates are used to host all readout
 197 electronics of the GlueX experiment. In addition to the VME-
 198 bus used to read out data from electronics modules, the VXS is
 199 instrumented with a high-speed serial bus in order to increase
 200 the bandwidth to several Gb/sec and provide an interconnected
 201 network between modules. The bus is used to transmit ampli-
 202 tudes digitized by the ADC to trigger electronics modules to
 203 include the CCAL in the Level 1 trigger system of the GlueX
 204 detector.

205 3.4. Light Monitoring System

206 To monitor performance of each calorimeter channel, we
 207 designed an LED-based light monitoring system (LMS). The
 208 LMS optics includes a blue LED, a spherical lens to correct
 209 the conical dispersion of the LED, and a diffusion grating to
 210 homogeneously mix the light. Light produced by the LED is
 211 incident on a bundle of plastic optical fibers (Edmund Optics)
 212 with a core diameter of 250 μm . Each fiber distributes light to
 213 an individual calorimeter module. On the crystal end, the fiber
 214 is attached to the module using a small acrylic cap glued to the
 215 crystal with a hole drilled through each cap to hold the fiber
 216 inside.

217 To monitor stability of the LED, we used two reference
 218 Hamamatsu 4125 PMTs, the same type as in the CCAL detec-
 219 tor. Each PMT receives light from two sources: a single fiber
 220 from the LED and a YAP:Ce pulser unit, both glued to the PMT
 221 face. The pulser unit consists of a 0.15 mm thick YAP:Ce scin-
 222 tillation crystal with a diameter of 3 mm spot activated by an
 223 ^{241}Am α source. The α source is used to monitor stability of
 224 the LED. The PMT was read out using a flash ADC. The high
 225 voltage on each reference PMT was adjusted to have the sig-
 226 nals from both the LED and α source fit within the range of a

12-bit flash ADC corresponding to 4096 counts, as shown in
 Fig. 5. Each LED was driven by a CAEN 1495 module, which
 allowed to generate LED pulses with a programmable rate. The
 LMS was integrated into the GlueX trigger system and provided
 a special trigger type during data taking. The LMS was exten-
 sively used during the detector commissioning and injected
 light to the CCAL detector with a typical frequency of 100 Hz
 continuously during the PrimEx- η experiment. This LED rate
 is similar to the trigger rate of events produced by the reference
 α source.

Most LMS components were positioned inside the
 temperature-stabilized detector box. The stability of the
 LED system measured using the reference PMTs during the
 entire PrimEx run was on the level of 1%. The ratio of signal
 ADC amplitudes from the LED pulser to the α source obtained
 during different run periods of the 48-day long PrimEx- η
 experiment is presented in Fig. 6. The ratio is normalized to
 the data in the beginning of the experiment. Stability of most
 CCAL modules observed using the LMS during the experiment
 was better than 6%. We did not apply any PMT gain adjust-
 ments during the experiment.

242 3.5. Calibration

243 The initial energy calibration of the CCAL was performed
 244 by moving the calorimeter platform and positioning each mod-
 245 ule into the photon beam during special low-intensity calibra-
 246 tion runs. The maximum rate in the module exposed to the
 247 beam did not exceed 200 kHz at a threshold of 15 MeV. The
 248 energy of each beam photon was determined by detecting a
 249 bremsstrahlung electron using the GlueX tagging detectors de-
 250 scribed in Section 2. The spot size of the collimated photon
 251 beam had a diameter of about 6 mm.

252 In the beginning of the calibration run, we adjusted the PMT
 253 high voltage for each module in order to equalize signal pulse
 254 amplitudes induced by 11 GeV beam photons. The amplitude
 255 was set to 3500 ADC counts, which corresponds to ~ 1.7 V.
 256 An example of flash ADC signal amplitude in the calorimeter

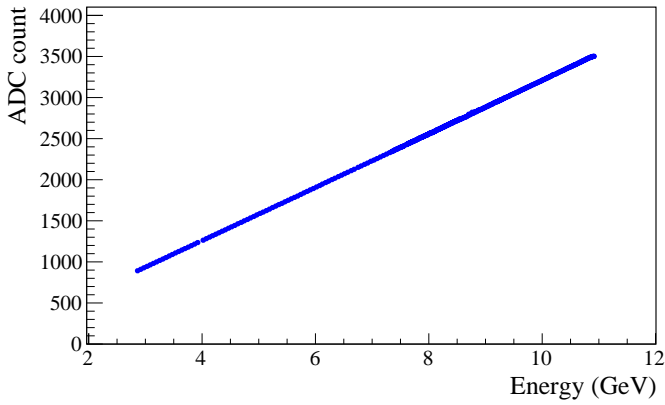


Figure 7: ADC signal pulse amplitude in the CCAL module as a function of the beam energy.

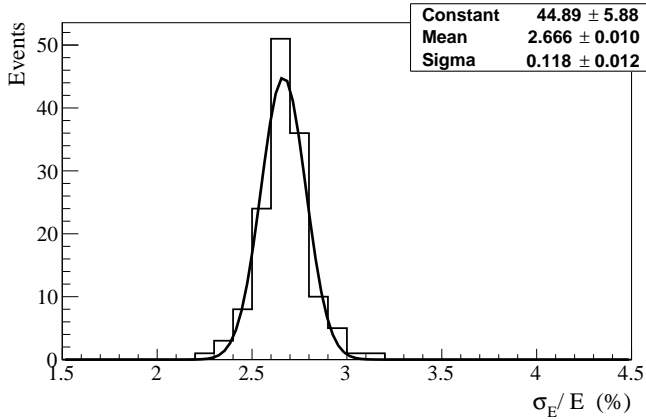


Figure 8: Relative energy resolution of 140 PbWO₄ modules installed on the CCAL measured with 6 GeV beam photons.

module as a function of the beam energy is presented in Fig. 7. The calibration of each module was refined by reconstructing showers in the calorimeter and constraining the reconstructed energy to the known beam energy.

During the calibration runs, we estimated the non-uniformity of the 140 CCAL modules by measuring the relative energy resolution for each individual module exposed to the beam. The energy resolution obtained for 6 GeV photons is presented in Fig. 8. The distribution is fit to a Gaussian function. The non-uniformity of the modules, i.e., the spread of the distribution is found to be smaller than 5%.

During calibration, we observed some non-linearity of the PMT active base with the large amplification factor of 24, on the level of a few percent, which impacted both the pulse peak and pulse integral. The base performance became linear when the amplifier gain was reduced. In order to study the impact of the non-linearity on the detector energy resolution, we replaced the original PMT active bases for 9 CCAL modules (in the array of 3 × 3 modules) with modified bases where the amplifier was bypassed. After adjusting high voltages and recalibrating PMT gains, we measured the energy resolution for different beam en-

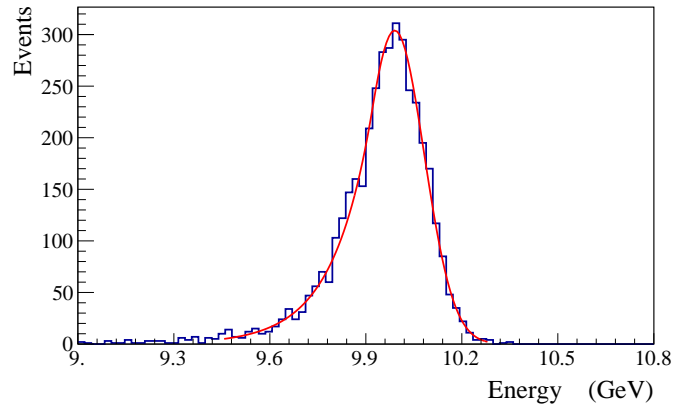


Figure 9: Energy distribution deposited by 10 GeV beam photons. The spectrum is fit to a Crystal Ball function.

ergies. The beam was incident on the center of the middle module in the array. An example of the energy deposited by 10 GeV photons is shown in Fig. 9. The energy resolution was obtained from a fit of the energy distribution to a Crystal Ball function³ implemented in the ROOT data analysis framework [17]. The energy resolution as a function of the beam energy is shown in Fig. 10. The distribution was fit to the following function:

$$\frac{\sigma_E}{E} = \frac{S}{\sqrt{E}} \oplus \frac{N}{E} \oplus C, \quad (1)$$

where S represents the stochastic term, N the electronic noise and C the constant term, E is the beam energy in GeV, and the symbol \oplus indicates a quadratic sum. The fit yields: $S = 2.63 \pm 0.01\%$, $N = 1.07 \pm 0.09\%$, and $C = 0.53 \pm 0.01\%$. The resolution was found to be about 10% better than that measured with the original base with the amplifier gain of 24⁴. The energy resolution is consistent with that of the HyCal calorimeter [13], which was instrumented with crystals produced by SICCAS in 2001 and was used in several experiments in Jefferson Lab's experimental Hall B. The HyCal PbWO₄ crystals have the same transverse size of 2.05 cm × 2.05 cm, but a smaller length of 18 cm. The initial CCAL calibration performed with the beam scan was fine-tuned after the PrimEx- η run by using showers of reconstructed Compton scattering candidates and constraining the reconstructed energy in the event to the know beam energy.

3.6. Performance during the PrimEx- η run

In the PrimEx- η experiment, we reconstruct Compton events produced by beam photons with the energy larger than 6 GeV. This energy range is covered by the GlueX pair spectrometer [18], which determines the photon flux needed for cross section measurements. An electron and photon produced in

³The function is named after the Crystal Ball collaboration.

⁴The linearity of the PMT active base is being currently improved; modified active bases will be installed before the new PrimEx- η run in 2021.

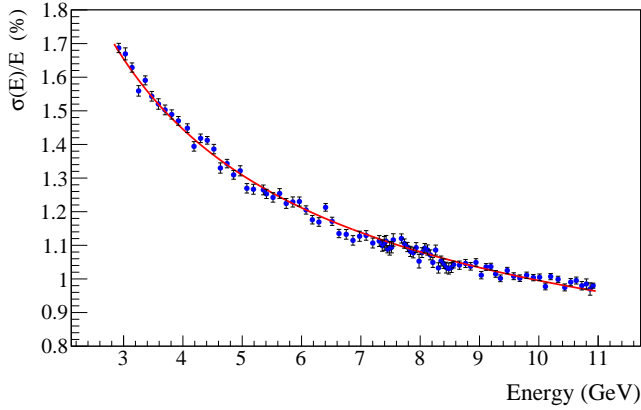


Figure 10: The CCAL energy resolution as a function of the photon energy.

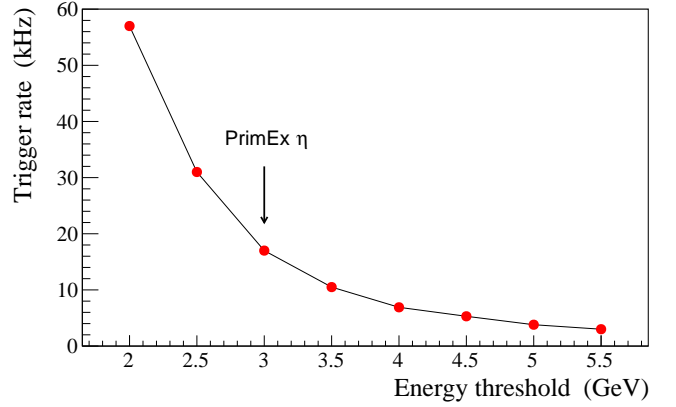


Figure 11: Trigger rate as a function of the total energy deposited in the FCAL and CCAL. The arrow indicates the energy threshold used in PrimEx- η production runs.

314 the Compton scattering process were detected by reconstruct-
 315 ing two showers, one in the FCAL and another one in the
 316 CCAL. The event topology of the reaction is such that the
 317 more energetic electron predominantly goes into the Compton
 318 calorimeter, while the photon strikes the FCAL. In order
 319 to accept Compton events during data taking and to reduce
 320 background originating from low-energy electromagnetic and
 321 hadronic interactions, the CCAL was integrated to the Level 1
 322 trigger system of the GlueX detector. The physics trigger was
 323 based on the total energy deposited in the forward and Compton
 324 calorimeters. The GlueX trigger is implemented on special-
 325 purpose programmable electronics modules with FPGA chips.
 326 The trigger architecture is described in Ref. [19]. The trigger
 327 rate as a function of the energy threshold is presented in Fig. 11.
 328 We collected data using a relatively small energy threshold of
 329 3 GeV at a trigger rate of about 18 kHz. This rate did not pro-
 330 duce any dead time in the data acquisition and trigger systems.
 331 The trigger rate was reproduced by a detailed Geant detector
 332 simulation.

333 The rate in the CCAL modules during the experiment is pre-
 334 sented in Fig. 12. In this plot, the photon beam goes through the
 335 center of the hole of 2×2 modules in the middle of the detector.
 336 The rate is the largest in innermost detector layers closest to the
 337 beamline. The maximum rate in the detector module was about
 338 200 kHz for an energy threshold of 30 MeV, which is equivalent
 339 to a signal pulse amplitude of 5 mV. Before the experiment, we
 340 performed a high-rate performance study of the PMT and elec-
 341 tronics using a laser and an LED pulser and did not find any
 342 degradation of the PMT gain up to 2 MHz [20].

343 Timing resolution of reconstructed showers is an important
 344 characteristic of the detector performance. In the experiment
 345 we used timing information provided by the calorimeters to
 346 identify the accelerator beam bunch for which the interaction
 347 occurred in the detector and therefore relate showers in the
 348 calorimeters with hits in the tagging detector, from the same
 349 event. A hit in the tagging detector defines the energy of the
 350 beam photon. The time in the calorimeter module is provided
 351 by an algorithm implemented on the programmable FPGA chip
 352 of the flash ADC. The algorithm performs a search of the peak

of the signal pulse and determines the time from the shape of
 the leading edge of the pulse. The times of all hits constituting
 the CCAL shower are combined to form the shower time by
 using an energy-weighted sum. The time difference between
 beam photon candidates and CCAL showers originating from
 Compton events is presented in Fig. 13. The main peak on this
 plot corresponds to beam photons and CCAL clusters produced
 in the same accelerator bunch. Satellite peaks, separated by the
 beam bunch period of about 4 ns, represent accidental beam
 photons from accelerator bunches not associated with the inter-
 action in the detector. The time resolution of CCAL showers
 is improved with the increase of the shower energy and was
 measured to be about 330 ps and 140 ps for 1 GeV and 9 GeV
 showers, respectively. In the PrimEx- η experiment the CCAL
 allowed a clear separation of beam photons originating from
 different beam bunches.

Reconstruction of electromagnetic showers in the FCAL is
 performed using an algorithm described in Ref. [21], which
 is a part of the standard GlueX reconstruction software. For
 the CCAL, we implemented an algorithm originally developed
 for the GAMS spectrometer [22, 23], which was subsequently
 adopted for the HyCal [13] in JLab's experimental Hall B. The
 algorithm provides a good separation of overlapping showers
 in the calorimeter by using profiles of electromagnetic showers.
 The elasticity distribution, defined as the reconstructed energy
 in the event minus the beam energy, is presented in Fig. 14 for
 Compton candidates produced by beam photons in the energy
 range between 6 GeV and 7 GeV. The solid line shows the fit
 of this distribution to the sum of a Gaussian and a second order
 polynomial function. The energy resolution of reconstructed
 Compton candidates in this energy range is about 130 MeV. In
 this plot, we subtracted background originating from accidental
 beam photons. This background was measured using off-time
 interactions and amounted to about 15%. The relatively small
 background, on the level of 10%, produced by interactions of
 the photon beam with the beamline material downstream the
 GlueX target was measured using empty-target runs and was
 also excluded from the elasticity distribution in Fig. 14. The

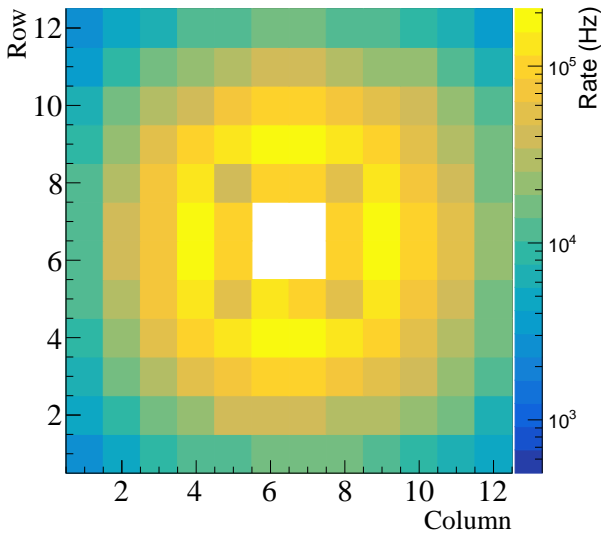


Figure 12: Rates in the CCAL modules during PrimEx- η production run. The energy threshold corresponds to 30 MeV. The beam goes through the center of the hole in the middle of the plot.

CCAL allowed to clearly reconstruct Compton events in the PrimEx- η experiment.

4. Upgrade of the GlueX forward calorimeter

The forward calorimeter of the GlueX detector consists of 2800 lead glass modules, each with a size of 4 cm \times 4 cm \times 45 cm, and is positioned about 6 m downstream of the target, as shown in Fig. 1. The FCAL covers a polar angle of photons produced from the target between 1 $^\circ$ and 11 $^\circ$ and detects showers with energies in the range of 0.1 - 8 GeV. The Cherenkov light produced in the module is detected by FEU-84-3 photomultiplier tubes, instrumented with Cockcroft-Walton bases [24]. The typical energy resolution of the FCAL is $\sigma_E/E = 6.2\%/\sqrt{E} \oplus 4.7\%$ [1].

The future physics program with the GlueX detector in Hall D will require an upgrade of the inner part of the forward calorimeter with high-granularity, high-resolution PbWO₄ crystals. The lead tungstate insert will improve the separation of clusters in the forward direction and the energy resolution of reconstructed photons by about a factor of two. Lead tungstate crystals possess better radiation hardness compared to lead glass, which is important for the long term operation of the detector at high luminosity. The size of the insert will tentatively comprise 1596 PbWO₄ crystals, which will form an array of 40 \times 40 modules⁵. Similar to the CCAL, the insert will have a beam hole of 2 \times 2 modules and a tungsten absorber used to shield the detector layer closest to the beamline. A schematic view of the FCAL frame with the installed lead tungstate insert

⁵The insert size proposed for the JEF experiment [2] is 1 m \times 1 m; the actual size will depend on the availability of funds.

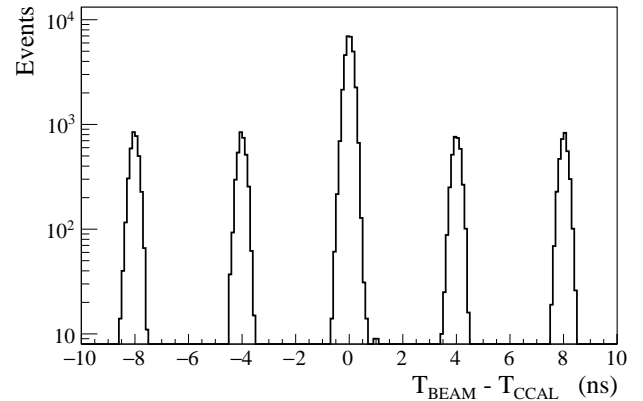


Figure 13: Time difference between beam photons and reconstructed CCAL showers for Compton candidates. Peaks are separated by the beam bunch period of 4 ns.

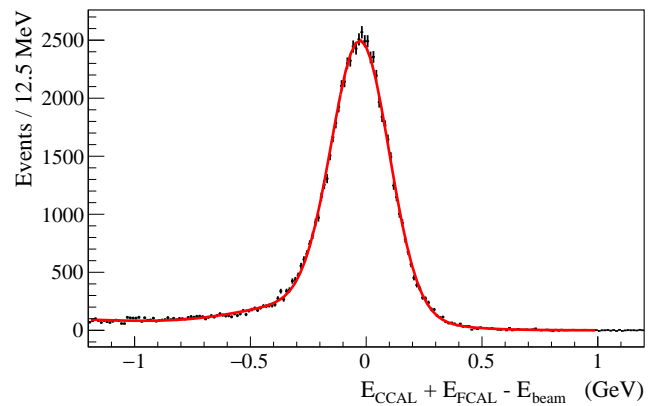


Figure 14: Elasticity distribution of reconstructed Compton candidates.

is presented in Fig. 15. Due to the different size of the lead glass bars and lead tungstate crystals, the lead glass modules stacked around the PbWO₄ insert will form four regions with a relative offset between modules; those regions are shown in green color in this plot.

The PbWO₄ module design of the FCAL insert will essentially be the same as for the CCAL, except for some small modifications needed to handle the magnetic field present in the FCAL region. The PMT housing made of the G-10 fiberglass material will be replaced by iron housing in order to reduce the magnetic field. The housing length will be increased to extend the magnetic shield beyond the PMT photocathode. An acrylic optical light guide will be inserted inside the PMT housing to couple the crystal and PMT.

The upgraded FCAL will be operated in GlueX experiments using a 30 cm long liquid hydrogen target at the designed photon flux of $5 \cdot 10^7$ γ /sec in the energy range between 8.4 GeV and 9 GeV. The designed luminosity is significantly larger than that used in the PrimEx- η experiment and was achieved after the PrimEx run in the fall of 2019. In order to finalize the design of the PMT electronics, it is important to understand detector rates

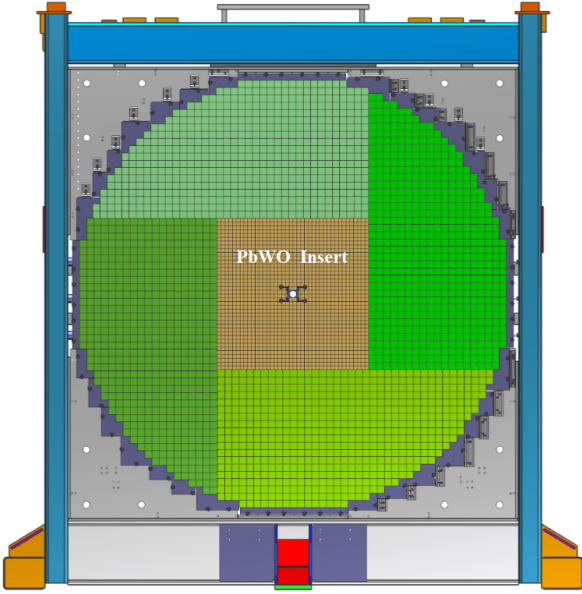


Figure 15: FCAL frame with calorimeter modules installed: PbWO₄ crystals (brown area), lead glass blocks (green). The photon beam passes through the hole in the middle of the calorimeter.

in the FCAL insert, especially in layers close to the beamline. We used the CCAL during high-intensity GlueX runs to study run conditions for the FCAL insert.

4.0.1. Magnetic shielding of PMTs

The longitudinal (directed along the beamline) and transverse (directed perpendicular to the axis of the beamline) components of the magnetic field produced by the GlueX solenoid magnet in the FCAL PbWO₄ insert area vary between 40 - 55 Gauss and 0 - 9 Gauss, respectively. The longitudinal field is the largest on the beamline, where the transverse component is practically absent. We studied the PMT magnetic shielding using a prototype consisting of an array of 3 × 3 PMT iron housings made of AISI 1020 steel, which was positioned in the middle of Helmholtz coils. Each housing had a size of 20.6 mm × 20.6 mm × 104 mm with a 20 mm round hole in the middle for the PMT. This corresponds to the realistic size of the magnetic shield that will be used in the calorimeter module assembly. Inside the housing we inserted two layers of mu-metal cylinders, with thicknesses of 350 μm and 50 μm, separated from each other by a Kapton film. The thickest cylinder was spot welded and annealed.

The Helmholtz coils had a diameter of about 1.5 m and can generate a uniform magnetic field with variable strength below 100 Gauss. A Hall probe was inserted into the central module of the prototype to measure the magnetic field at different Z-positions along the length of the cylinder. The field was measured for two different orientations of the prototype with respect to the magnetic field: field oriented along the PMT (longitudinal, B_z) and perpendicular to the PMT housing (transverse, B_x). Field measurements are presented in Fig. 16. The PMT shield significantly reduces both the longitudinal and transverse fields to the level of $B_z \sim 1$ Gauss and $B_x \ll 1$ Gauss. The

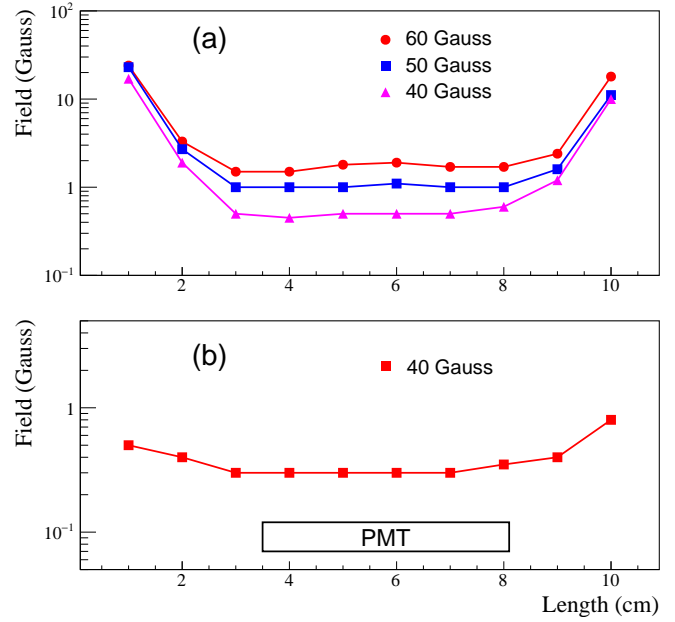


Figure 16: Magnetic field distribution inside the PMT shield housing as a function of the distance from the housing face. Plot (a) corresponds to the longitudinal field and plot (b) corresponds to the transverse field produced by the Helmholtz coils. Markers denote different field values.

transverse field, which is well shielded, is more critical for the PMT operation, as it is directed perpendicular to the electron trajectory inside the photo tube and deflects electrons, resulting in the degradation of the photon detector efficiency and gain. The field reaches a plateau at $Z = 3$ cm from the face of the housing. We will use 3.5 cm long acrylic light guides, in order to place the most sensitive to the magnetic field area of the PMT between the photocathode and the last dynode (4.6 cm long) in the region with the smallest magnetic field, as shown in Fig. 16. The actual field inside the FCAL insert module is expected to be even smaller due to the collective shielding effect, i.e., the large amount of shielding material installed on surrounding modules [25].

We studied performance of the shielded PMT in the magnetic field using an LED pulser. A blue LED with a light diffuser was placed about 20 cm from the PMT housing prototype and was aligned with the middle module. The PMT response was measured for different pulse amplitudes and operational high voltages. In order to study the contributions from longitudinal and transverse field components we rotated the prototype by different angles. Signal amplitudes as a function of the magnetic field measured in the prototype tilted by about 10 degrees are presented on the top plot of Fig. 17. Amplitudes, normalized to measurements without magnetic field, are shown on the bottom plot. The relative degradation of the signal amplitude for the maximum field in the FCAL insert region of $B = 55$ Gauss ($B_z \sim 54$ Gauss and $B_x \sim 9$ Gauss) was measured to be on the level of 1%. The proposed shielding configuration is sufficient to reduce the magnetic field to the level suitable for the PMT operation.

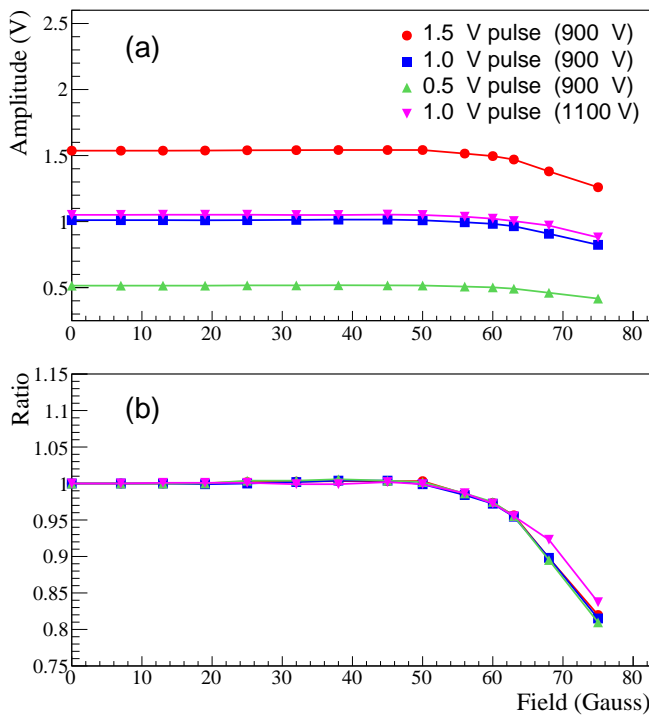


Figure 17: Signal amplitudes of shielded PMT induced by an LED as a function of the magnetic field (a). Amplitudes, normalized to measurements without magnetic field (b). The PMT response was measured for different intensities of light pulse and HV settings as shown by different polymarkers.

4.0.2. Light guide studies

Studies of the magnetic shielding demonstrated that the PMT has to be positioned inside the iron housing at the distance of at least 3 cm from the face of the PbWO_4 crystal. In order to do this, in the FCAL insert module we decided to use a 3.5 cm long acrylic cylindrical light guide with a diameter of 18.5 mm between the PMT and the PbWO_4 crystal. The light guide is wrapped with reflective ESR foil and attached to the PMT with Dymax 3094 UV curing glue. Optical coupling to the crystal is provided by a “silicon cookie”: a 1 mm thick transparent rubber cylinder made of the room temperature vulcanized silicon compound, RTV615. The silicon cookie is not glued to the light guide or the crystal, so the module can be easily disassembled if its PMT needs to be replaced.

We compared light losses of the FCAL insert module instrumented with the light guide with the CCAL module, where the PMT was coupled directly to the crystal using an optical grease. Light collection was measured using electrons provided by the Hall D pair spectrometer (PS) [18]. The PS is used to measure the flux of beam photons delivered to the experimental hall by detecting electromagnetic electron-positron pairs produced by the photons in a thin converter inserted to the beam. Leptons from the pair are deflected in a dipole magnet and registered using scintillator detectors placed in the electron and positron arms of the spectrometer. The energy of a lepton is detected using a high-granularity PS hodoscope, which consists of 145

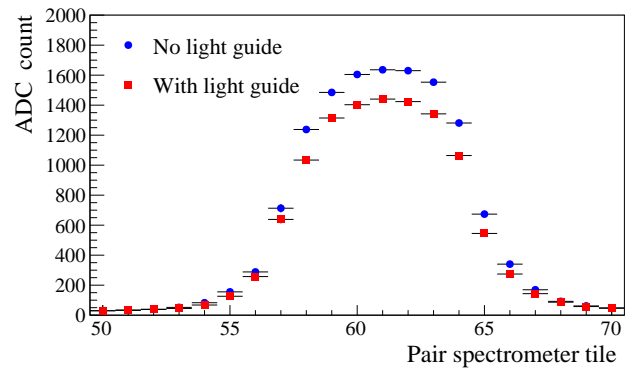


Figure 18: ADC amplitudes of the calorimeter module as a function of the pair spectrometer tile for two configurations: the PMT directly coupled to the PbWO_4 crystal (circles), and the PMT coupled to the module using optical light guide (boxes).

scintillating tiles and covers the energy range between 3 GeV and 6 GeV. Each tile corresponds to the specific lepton energy.

The relative light yield of the module with and without the light guide was estimated by positioning the module behind the PS and measuring signal amplitudes induced by the PS electrons. We first measured the ADC response in the CCAL module, which was subsequently modified by adding the light guide to the same PMT and crystal and was placed to the same spot of the PS test setup. Results of the measurements are presented in Fig. 18. The ADC amplitude of the calorimeter module is presented as a function of the PS tile for the two module configurations with and without the light guide. The light guide results in a relatively small loss of light of 15 – 20% compared with the CCAL module. We note that wrapping the light guide with the reflective material is important. Losses in unwrapped light guide constitute about 35%. We repeated light collection measurements using two more modules and obtained consistent results.

4.0.3. Detector rate

The PMT anode current is one of the critical characteristics that have to be considered during the design of the PMT divider. Typically the anode current should be on the level of a few micro amperes and significantly smaller than the divider current in order to provide stable performance of the PMT base and prevent the long-term degradation of the PMT. Some lifetime tests of the Hamamatsu 4125 PMT are described in Ref. [26].

The anode current (I) was measured in the CCAL modules during data production runs at the GlueX nominal luminosity. It was obtained by measuring the average voltage in the flash ADC induced by particles incident on the CCAL module as follows:

$$I = \frac{\bar{U}}{R} \cdot \frac{1}{G}, \quad (2)$$

where \bar{U} is the average voltage in units of Volts, R is the input impedance of the amplifier ($\sim 50 \Omega$), and G is the amplifier

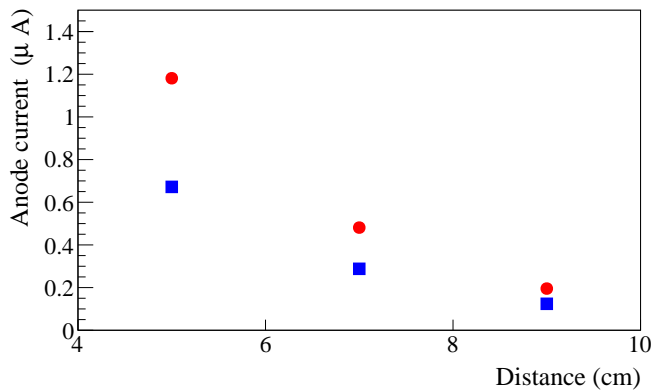


Figure 19: Typical PMT anode current of CCAL modules positioned at different distances from the beamline. Circles correspond to the nominal GlueX luminosity, boxes correspond to 60% of the nominal luminosity.

gain of 24. A periodic pulser not associated with an interaction in the detector was used as a trigger to read out flash ADC raw data for each CCAL module in a time window of 400 ns. The voltage was determined by summing up ADC amplitudes in the readout window and normalizing the sum to the window size. The typical anode current measured in CCAL modules situated at different distances from the beamline is presented in Fig. 19. Modules from the first CCAL layer closest to the beamline and the outer most layer were not used in the analysis. The inner modules were shielded by a tungsten absorber and the outer modules were obscured by the FCAL. The rate in the detector is dominated by the forward-directed electromagnetic background. The estimated anode current is the largest in the innermost layer of the detector closest to the beamline and amounts to about $1.4 \mu\text{A}$. This current is significantly smaller than the PMT divider current of about $300 \mu\text{A}$.

We used the CCAL measurements to estimate the current in the FCAL insert. Taking the geometrical location of FCAL and CCAL modules into account, the largest PMT current in the FCAL insert modules closest to the beamline and not shielded by the absorber was conservatively estimated to be about $15 \mu\text{A}$. We assume that the PMT base is operated at 1 kV and no amplifier is used. The detector rate drops rapidly with the increase of the radial distance from the beamline. The estimated anode current is relatively large and must be reduced by lowering the PMT high voltage. We are considering to instrument PMTs in a few inner FCAL insert layers with an amplifier with a gain of 5 and to omit the amplifier on other modules. We are planning to perform more beam tests of the FCAL insert active base using the CCAL in forthcoming GlueX runs in 2021-2022.

5. Neutral Particle Spectrometer

The NPS is a new facility in Hall C that will allow access to precision measurements of small cross sections of reactions with neutral final states. The NPS consists of an electromag-

netic calorimeter preceded by a sweeping magnet. As operated in Hall C, it replaces one of the focusing spectrometers.

The NPS science program currently features six fully approved experiments. E12-13-010 [4] and E12-06-114 [5] experiments will measure the Exclusive Deeply Virtual Compton Scattering and π^0 cross sections to the highest Q^2 accessible at Jefferson Lab. Both experiments will provide important information for understanding Generalized Parton Distributions (GPDs). The E12-13-007 [6] experiment will study semi-inclusive π^0 electroproduction process and seeks to validate the factorization framework that is needed by the entire 12 GeV Jefferson Lab semi-inclusive deep-inelastic scattering program. Measurements of Wide-Angle and Timelike Compton Scattering reactions will be performed by the E12-14-003 [7] and E12-17-008 [8] experiments. These measurements will allow to test universality of GPDs using high-energy photon beams. The NPS will also be used in the E12-14-005 [9] experiment to study exclusive production of π^0 at large momentum transfers in the process $\gamma p \rightarrow \pi^0 p$.

The NPS science program requires neutral particle detection over an angular range between 6 and 57.3 degrees at distances of between 3 and 11 meters from the experimental target. The experiments will use a high-intensity beam of electrons with the energies of 6.6, 8.8, and 11 GeV, and a typical luminosity of $\sim 10^{38} \text{ cm}^{-2}\text{s}^{-1}$ as well as a secondary beam of photons incident on a liquid hydrogen target. A vertical-bend sweeping magnet with integrated field strength of 0.3 Tm will be installed in front of the spectrometer in order to suppress and eliminate background of charged particle tracks originating from the target. The photon detection is the limiting factor of the experiments. Exclusivity of the reaction is ensured by the missing mass technique and the missing-mass resolution is dominated by the energy resolution of the calorimeter. The calorimeter is anticipated to provide the spacial resolution of 2-3 mm and the energy resolution of about $2.5\% / \sqrt{E}$. The NPS consists of 1080 PbWO_4 crystals that form an array of 30×36 modules. Similarly to the FCAL insert in Hall D, the NPS will be built from the crystals of the same size, and instrumented with the same type of PMTs and readout electronics. The details of the mechanical assembly and commissioning of the NPS are currently under development and will be described in a forthcoming publication.

The radiation hardness and good optical quality of lead tungstate crystals are critical for the NPS calorimeter. The NPS collaboration, in a synergistic effort with the EIC eRD1 consortium, has characterized to date over 1200 PbWO_4 crystals produced by CRYTUR and SICCAS from 2014 to the present. The results of these studies have been published in Ref. [10]. CRYTUR crystal samples were found to have greater overall uniformity in transmittance and light yield, and better radiation hardness. Of the samples characterized by the NPS collaboration 140 SICCAS crystals have been used in the CCAL detector.

⁶The minimum NPS angle at 3 m is 8.5 degrees; at 4 m it is 6 degrees.

6. Summary

We described the design and performance of the Compton calorimeter, which was constructed using 140 lead tungstate PbWO_4 crystals recently produced by SICCAS. The calorimeter was successfully used in the PrimEx- η experiment in spring of 2019 for reconstruction of Compton scattering events. The CCAL served as a prototype for two large-scale electromagnetic calorimeters based on the PbWO_4 crystals: the lead tungstate insert of the forward calorimeter of the GlueX detector and the neutral particle spectrometer. Experience gained during construction and operation of the CCAL provided important information for finalizing the design of FCAL PbWO_4 modules and PMT dividers and also served to further optimize the NPS calorimeter. We presented the design of the FCAL lead tungstate insert and gave an overview of the NPS project.

7. Acknowledgments

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