

# $\rho(770)$ Meson Spin-Density Matrix Elements

Discussion of Uncertainties

Alexander Austregesilo

Curtis A. Meyer

Naomi S. Jarvis

Amplitude Analysis WG Meeting

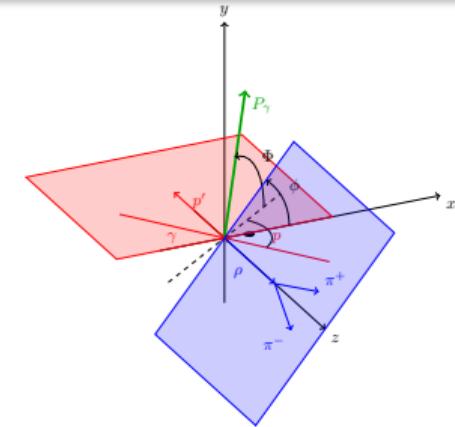
October 17, 2022



# Production Mechanism

## Spin-Density Matrix Elements

- Full angular distribution of vector meson production and decay is described by **spin-density matrix elements**  $\rho_{ij}^k$
- Linear beam polarization provides access to **nine** linearly independent SDMEs
- Intensity **W** is expressed as function of angles  **$\cos \vartheta, \varphi, \Phi$**  and degree of polarization  **$P_\gamma$**



$$W(\cos \vartheta, \varphi, \Phi) = W^0(\cos \vartheta, \varphi) - P_\gamma \cos(2\Phi) W^1(\cos \vartheta, \varphi) - P_\gamma \sin(2\Phi) W^2(\cos \vartheta, \varphi)$$

$$W^0(\cos \vartheta, \varphi) = \frac{3}{4\pi} \left( \frac{1}{2}(1 - \rho_{00}^0) + \frac{1}{2}(3\rho_{00}^0 - 1) \cos^2 \vartheta - \sqrt{2}\text{Re}\rho_{10}^0 \sin 2\vartheta \cos \varphi - \rho_{1-1}^0 \sin^2 \vartheta \cos 2\varphi \right)$$

$$W^1(\cos \vartheta, \varphi) = \frac{3}{4\pi} \left( \rho_{11}^1 \sin^2 \vartheta + \rho_{00}^1 \cos^2 \vartheta - \sqrt{2}\text{Re}\rho_{10}^1 \sin 2\vartheta \cos \varphi - \rho_{1-1}^1 \sin^2 \vartheta \cos 2\varphi \right)$$

$$W^2(\cos \vartheta, \varphi) = \frac{3}{4\pi} \left( \sqrt{2}\text{Im}\rho_{10}^2 \sin 2\vartheta \sin \varphi + \text{Im}\rho_{1-1}^2 \sin^2 \vartheta \sin 2\varphi \right)$$

Schilling *et al.* [Nucl. Phys. B, 15 (1970) 397]

$$W(\cos \vartheta, \varphi, \Phi) = W^0(\cos \vartheta, \varphi) - P_\gamma \cos(2\Phi) W^1(\cos \vartheta, \varphi) - P_\gamma \sin(2\Phi) W^2(\cos \vartheta, \varphi)$$

Measured Intensity  $I(\Omega) \propto W(\cos \vartheta, \varphi, \Phi)$

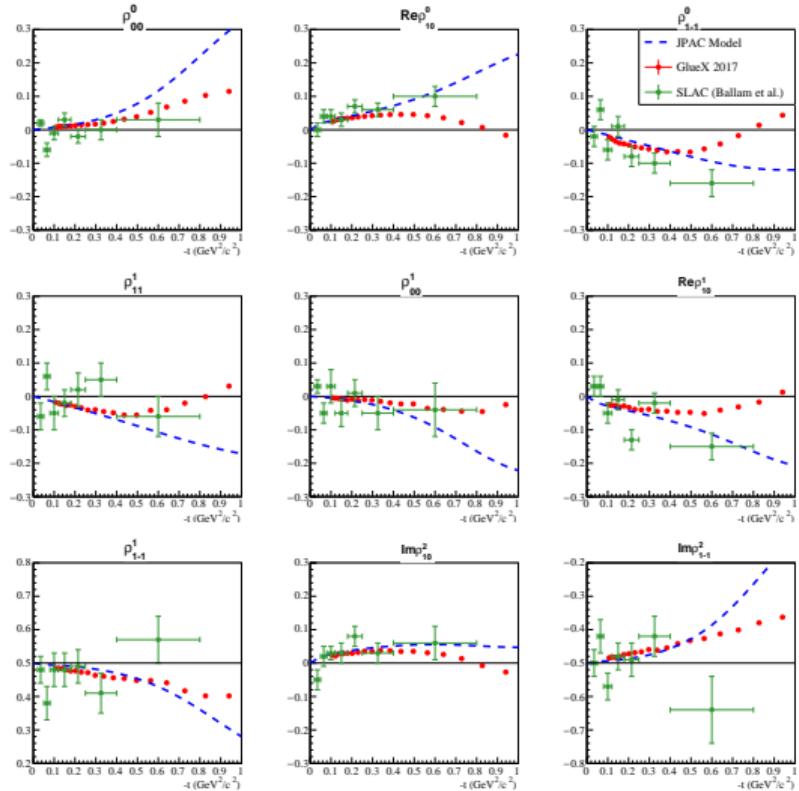
## Extended Maximum-Likelihood Fit

$$\ln L = \underbrace{\sum_{i=1}^N \ln I(\Omega_i)}_{\text{Signal Events}} - \underbrace{\sum_{j=1}^M \ln I(\Omega_j)}_{\text{Background}} - \underbrace{\int d\Omega I(\Omega) \eta(\Omega)}_{\text{Normalization Integral}}$$

- Maximize by choosing SDMEs such that the intensity fits the observed  $N$  events
- Accidental background subtracted in likelihood
- Normalization integral evaluated by a phase-space Monte Carlo sample with the acceptance  $\eta(\Omega) = 0/1$

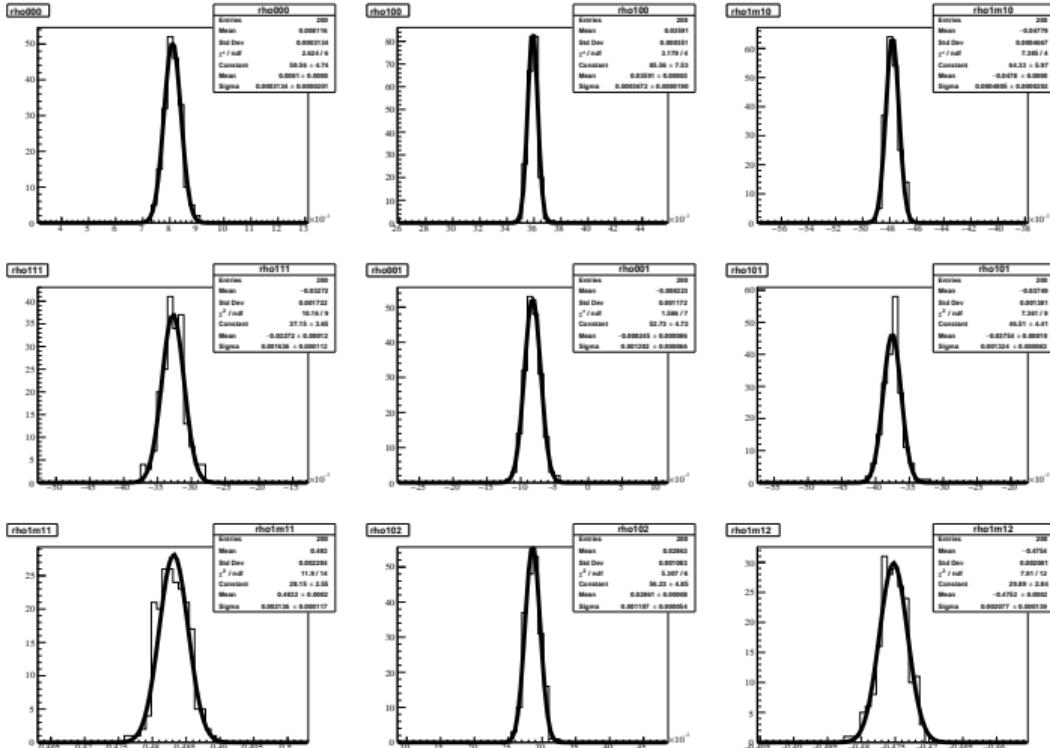
# Result

$\gamma p \rightarrow \rho(770)p$



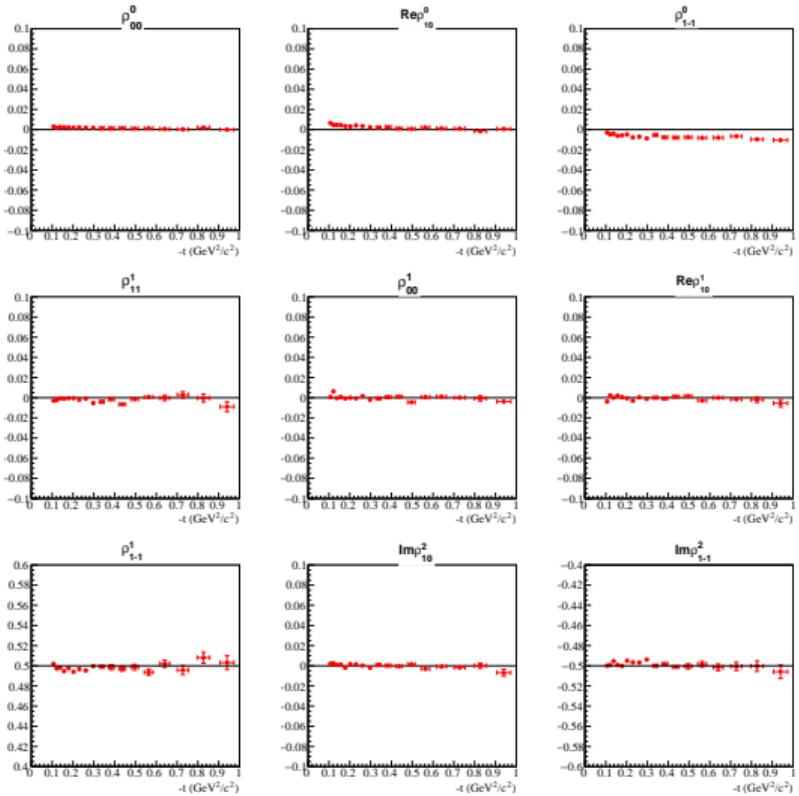
- Combined fit of 4 orientations with constraints
- Excellent agreement with JPAC for  $t < 0.5 \text{ GeV}^2$
- Statistical uncertainties only
- Systematic studies presented today

# Statistical Uncertainties: Bootstrapping



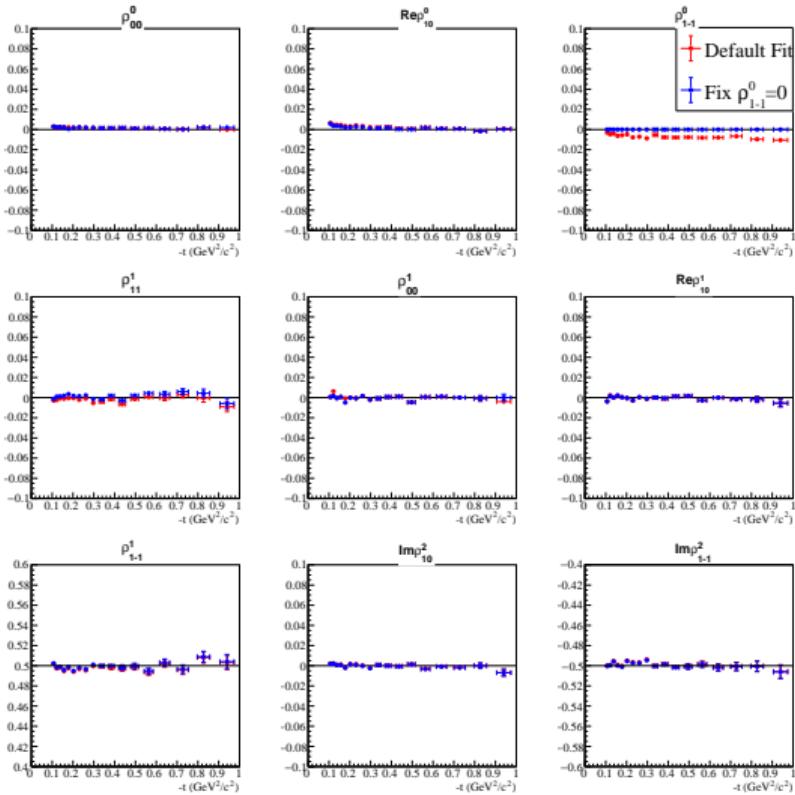
- Repeat fit 200 times by resampling the datasets
- Determine mean and variance via Gaussian fit, use for final result
- Gaussian variance  $\approx 25\%$  larger than Minuit uncertainty

# Input/Output Test with Signal MC



- Significant effect for  $\rho_{1-1}^0$  for full  $t$  range
- Smaller effect on  $\rho_{10}^0$  near  $t = 0.1$

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- Significant effect for  $\rho_{1-1}^0$  for full  $t$  range
- Smaller effect on  $\rho_{10}^0$  near  $t = 0.1$
- Nearly independent from other SDMEs
- Add deviation to systematic uncertainty

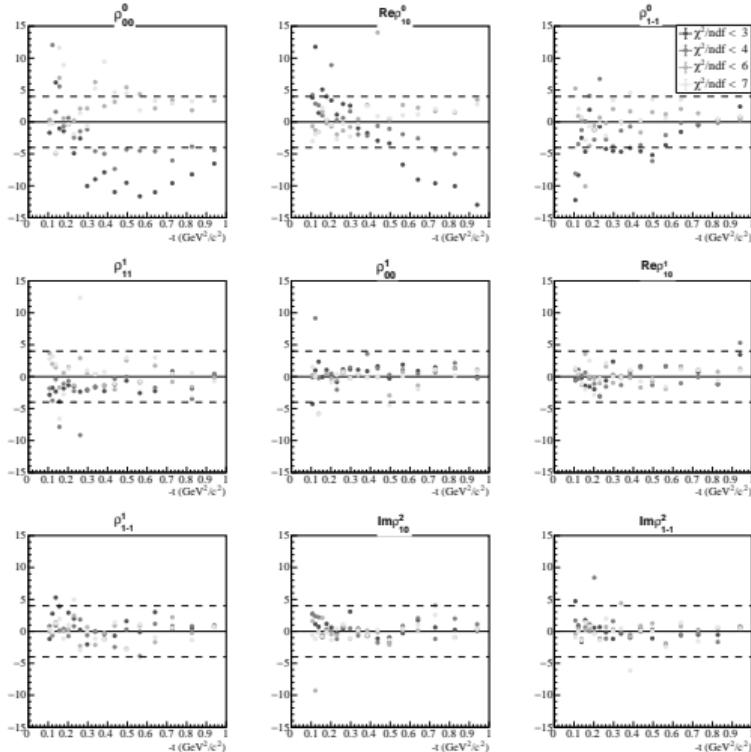
# Barlow's Significance Test

$$B = \frac{\Delta}{\sigma_B} = \frac{\rho - \rho_i}{\sqrt{|\sigma^2 - \sigma_i^2|}}$$

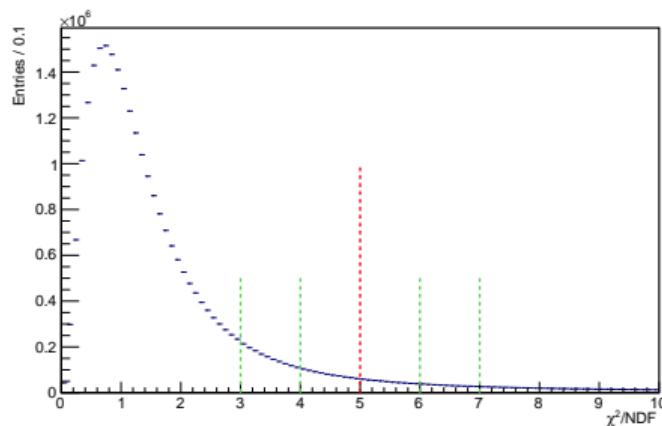
- Gauge statistical significance suggested by Barlow
- If  $B < 1$ : not significant
- If  $B > 4$ : take into account
- Else: discuss

arXiv:hep-ex/0207026

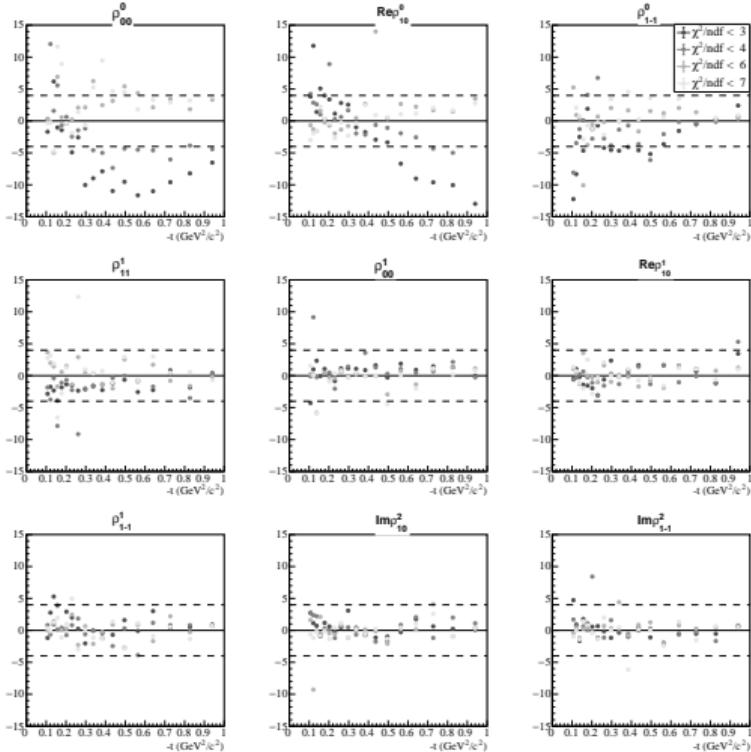
# Kinematic Fit



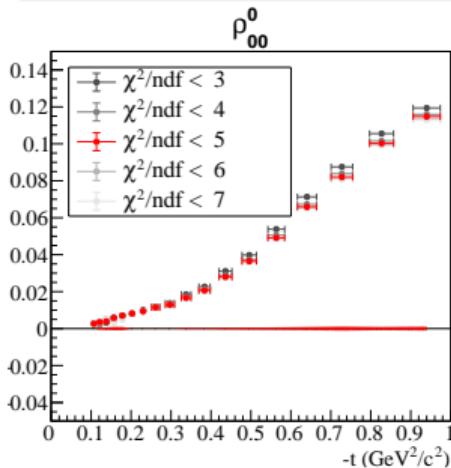
- Default:  $\chi^2/\text{ndf} < 5$
- Variation:  $\pm 2$  corresponds roughly to  $\pm 10\%$  data
- Unpolarized  $\rho^0$ 's clearly fail significance test



# Kinematic Fit

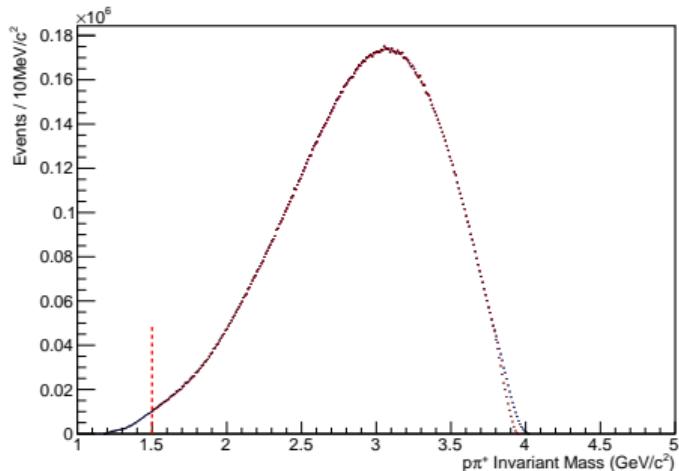


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- Compute standard deviation from variations and use as systematic uncertainty

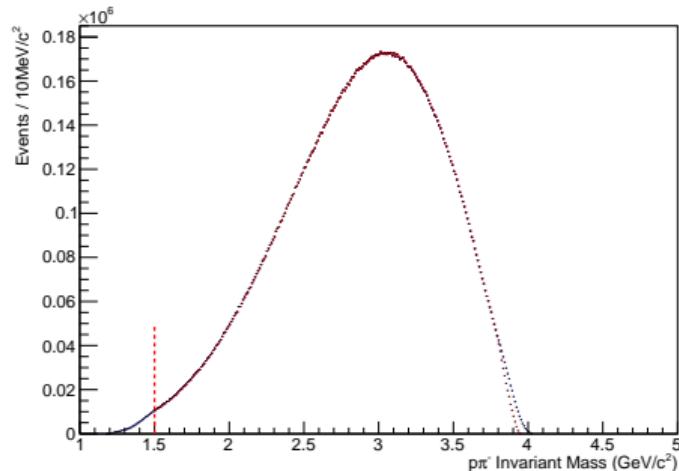


# Effect of Target Excitation

$p\pi^+$  Invariant Mass

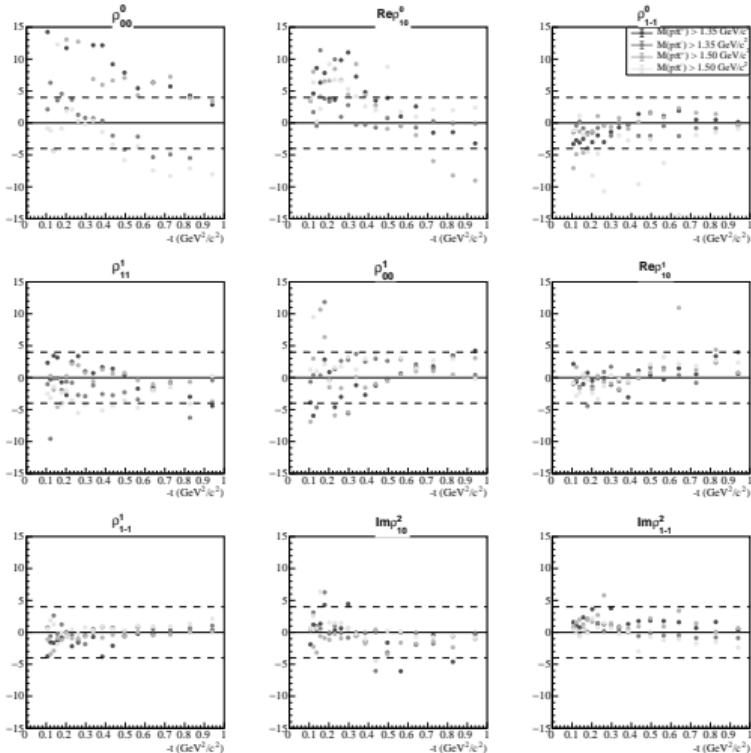


$p\pi^-$  Invariant Mass

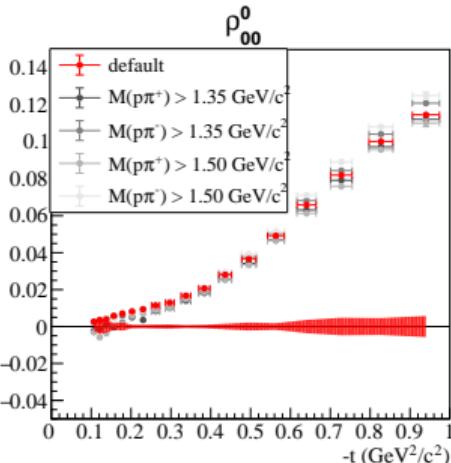


- Nearly no evidence for baryon excitations after selection of  $\rho(770)$  mass region
- Systematic study: cuts at  $M(p\pi^\pm) > 1.35, 1.5 \text{ GeV}/c^2$ , data reduction maximal 0.6%

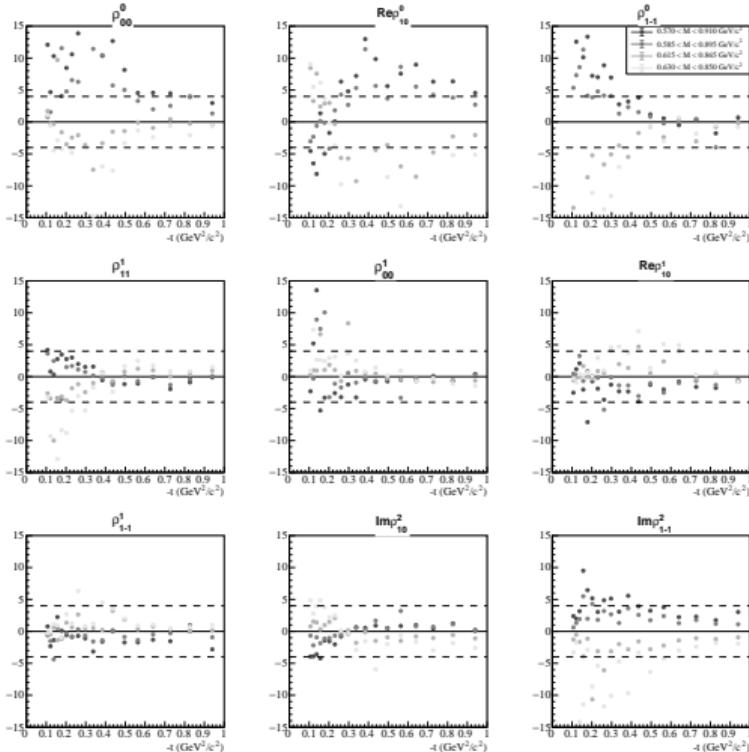
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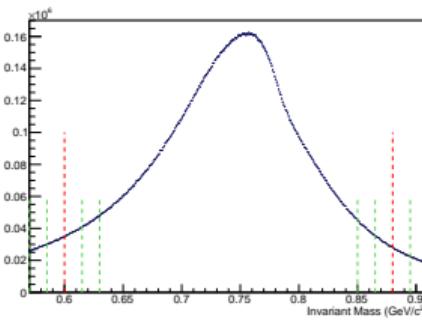
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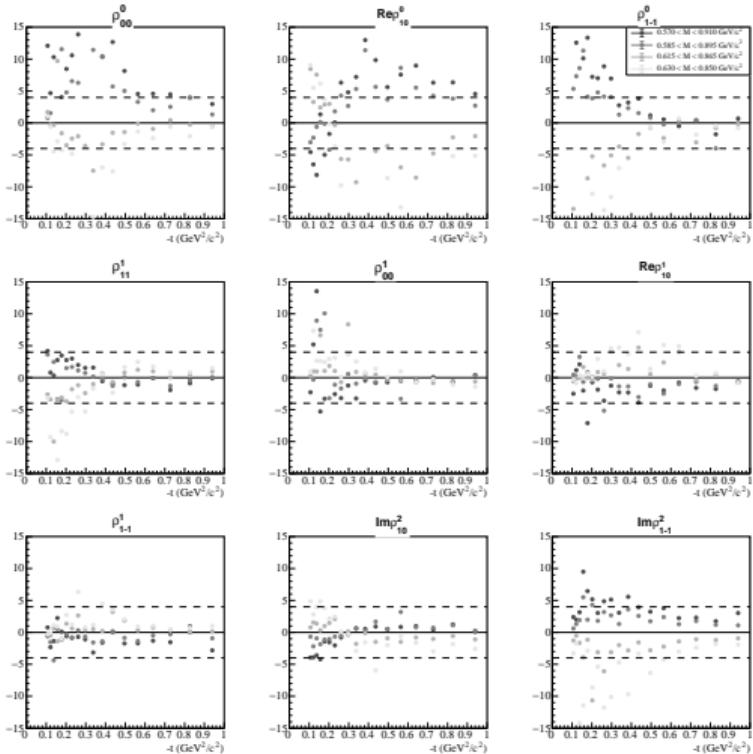
# Invariant Mass



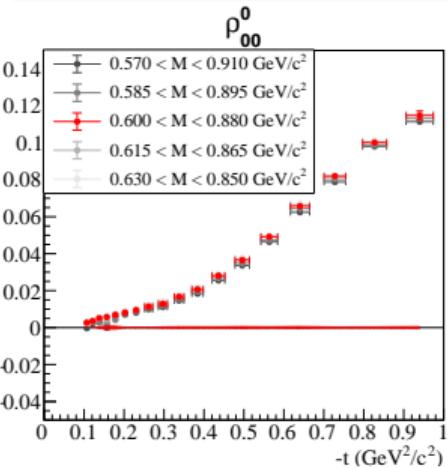
- Variation:  $\pm 150, 300 \text{MeV}/c^2$  corresponds to maximum of  $\pm 10\%$  data
- Several SDMEs fail significance test
- Compute standard deviation from variations and use as systematic uncertainty



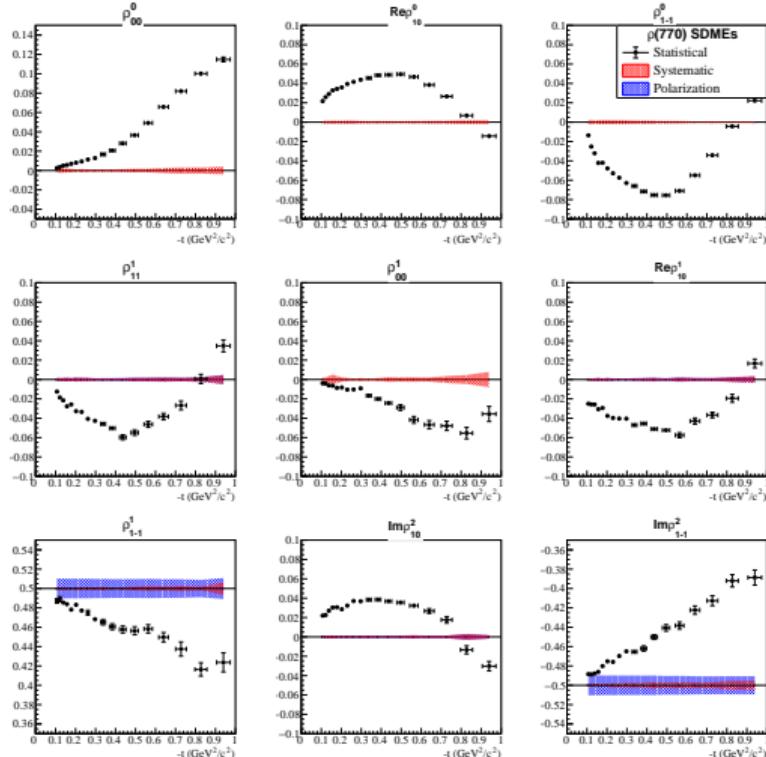
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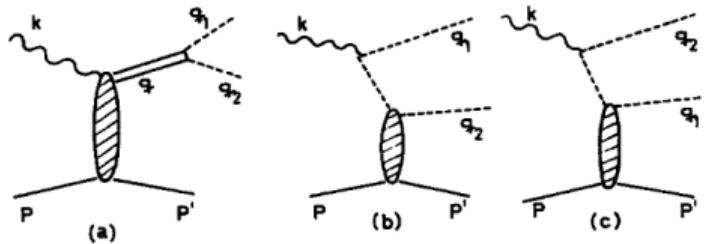


# Beam Polarization and Summary

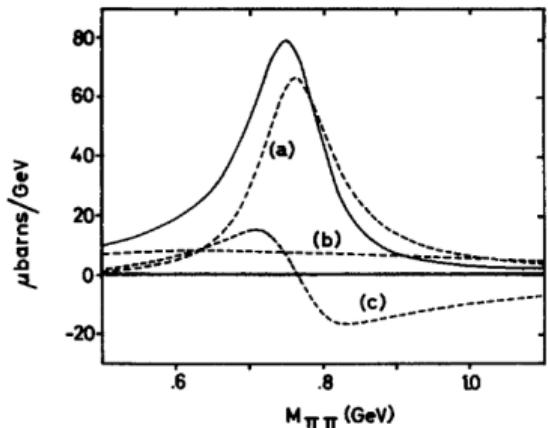


- Individual contributions added quadratically (plot has to be checked)
- 2.1% systematic uncertainty on  $P_\gamma$  added quadratically to  $\rho^{1,2}$ 's only
- Uncertainty from Input/Output test not yet added
- No significant contribution from orientations of polarization

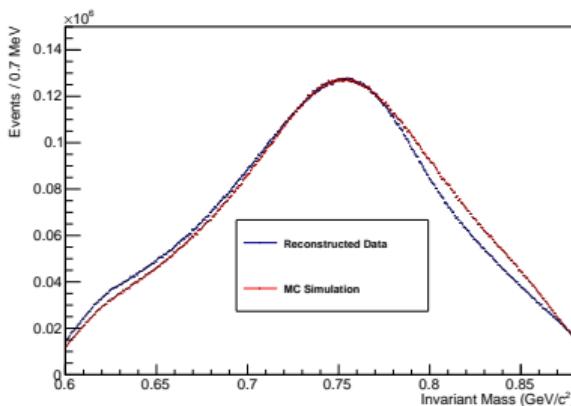
# Excursion: Mass Dependence



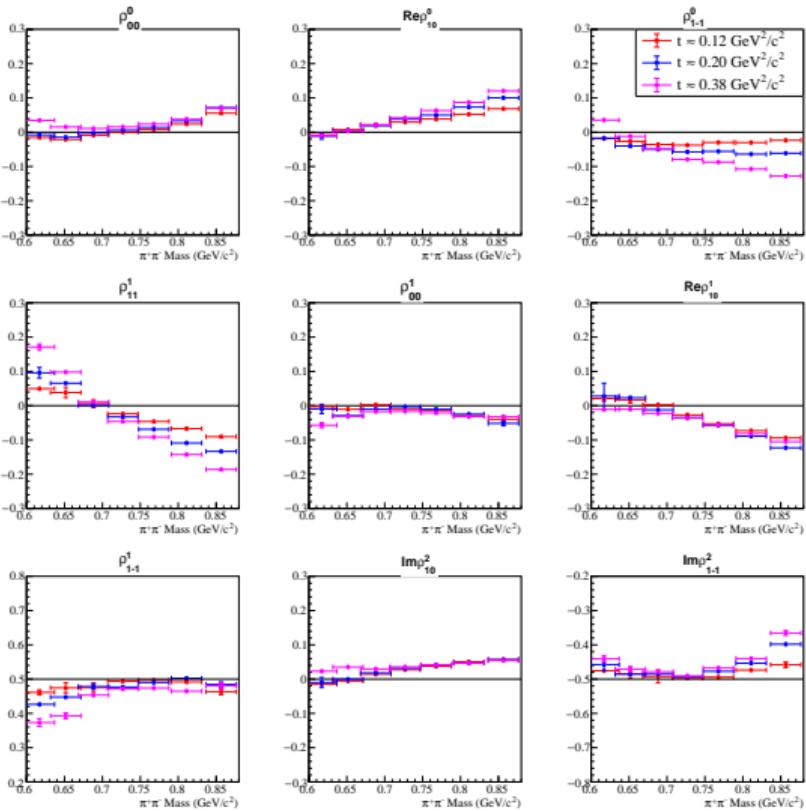
- Breit-Wigner Mass observed to be  $18 \text{ MeV}/c^2$  lower than PDG value
- Consistent with earlier observations
- Explained by interference with non-resonant processes (b) and (c)



Soding, Phys. Lett. 19, 702 (1966)



# Excursion: Mass Dependence



- Extract SDMEs as a function of  $\pi^+\pi^-$  mass for each  $t$  bin
- Mass dependence increases with  $t$  as S-wave background increases
- $\rho_{11}^1$  shows largest mass dependence, but effect seen in all
- Physics result instead of systematic error
- Narrower mass bin will increase statistical uncertainty

# Excursion: Polarized Reflectivity Amplitudes



M.R. Shepherd, [GlueX doc 4094]

- Definition:

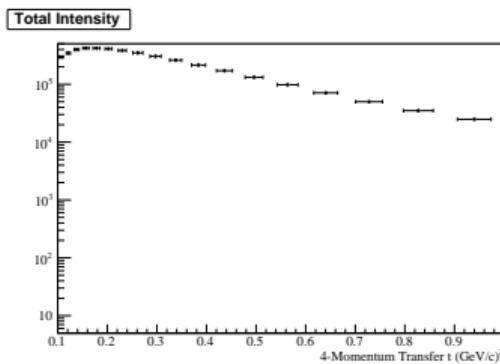
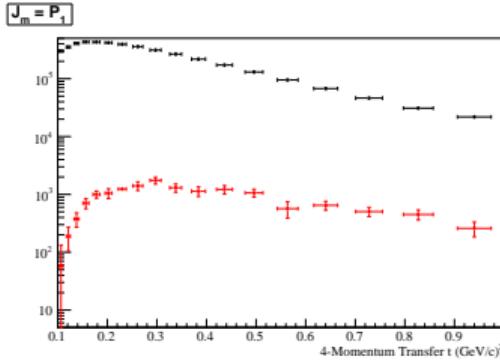
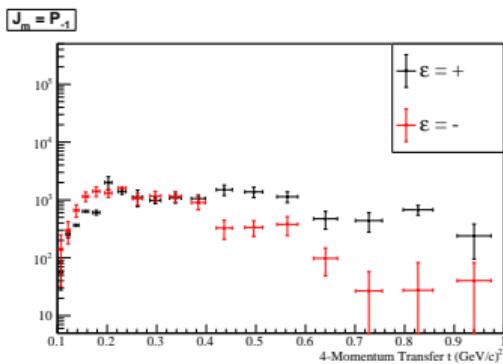
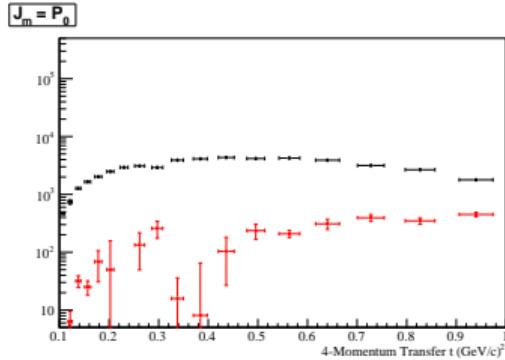
$$Z_\ell^m(\Omega, \Phi) \equiv Y_\ell^m(\Omega) e^{-i\Phi}.$$

- Final formulation of intensity:

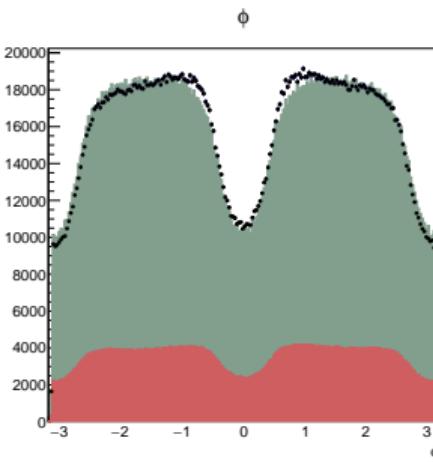
$$I(\Omega, \Phi) = 2\kappa \sum_k \left\{ (1 - P_\gamma) \left| \sum_{\ell,m} [\ell]_{m;k}^{(-)} \operatorname{Re}[Z_\ell^m(\Omega, \Phi)] \right|^2 + (1 - P_\gamma) \left| \sum_{\ell,m} [\ell]_{m;k}^{(+)} \operatorname{Im}[Z_\ell^m(\Omega, \Phi)] \right|^2 + (1 + P_\gamma) \left| \sum_{\ell,m} [\ell]_{m;k}^{(+)} \operatorname{Re}[Z_\ell^m(\Omega, \Phi)] \right|^2 + (1 + P_\gamma) \left| \sum_{\ell,m} [\ell]_{m;k}^{(-)} \operatorname{Im}[Z_\ell^m(\Omega, \Phi)] \right|^2 \right\}.$$

- Absorb factor  $\sqrt{1 \pm P_\gamma}$  into amplitude
- Repeated  $[\ell]_{m;k}^\pm$  have to be constrained
- Equivalent to decomposition into SDMEs

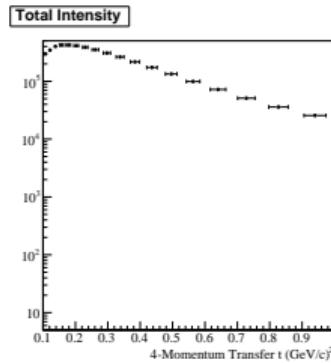
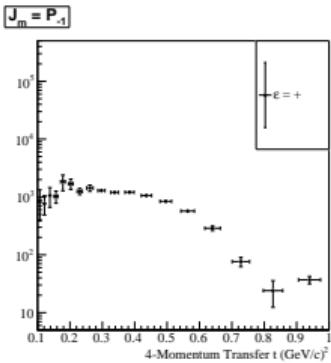
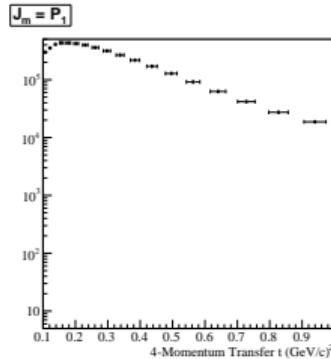
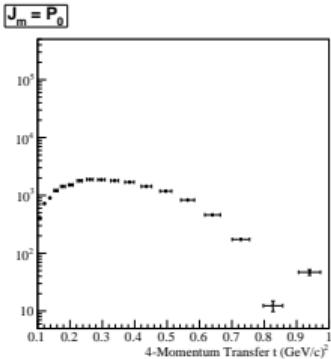
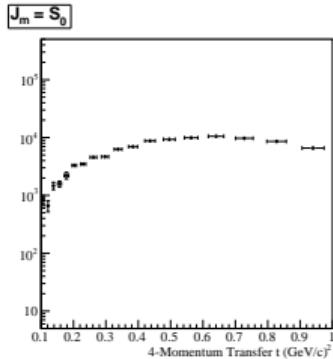
# Excursion: Polarized Reflectivity Amplitudes



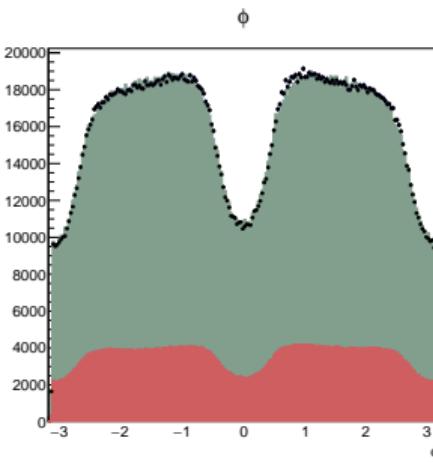
- Stable fit with all  $P$ -waves:  
 $P^{\varepsilon=\pm}_{\ell=-1,0,1}$
- Small negative reflectivity contribution



# Excursion: Polarized Reflectivity Amplitudes



- Stable fit with all  $P$ -waves:  
 $P^{\varepsilon=\pm}$   
 $\ell=-1,0,1$
- Small negative reflectivity contribution neglected
- Interference with  $S$ -wave explains  $\varphi$  asymmetry



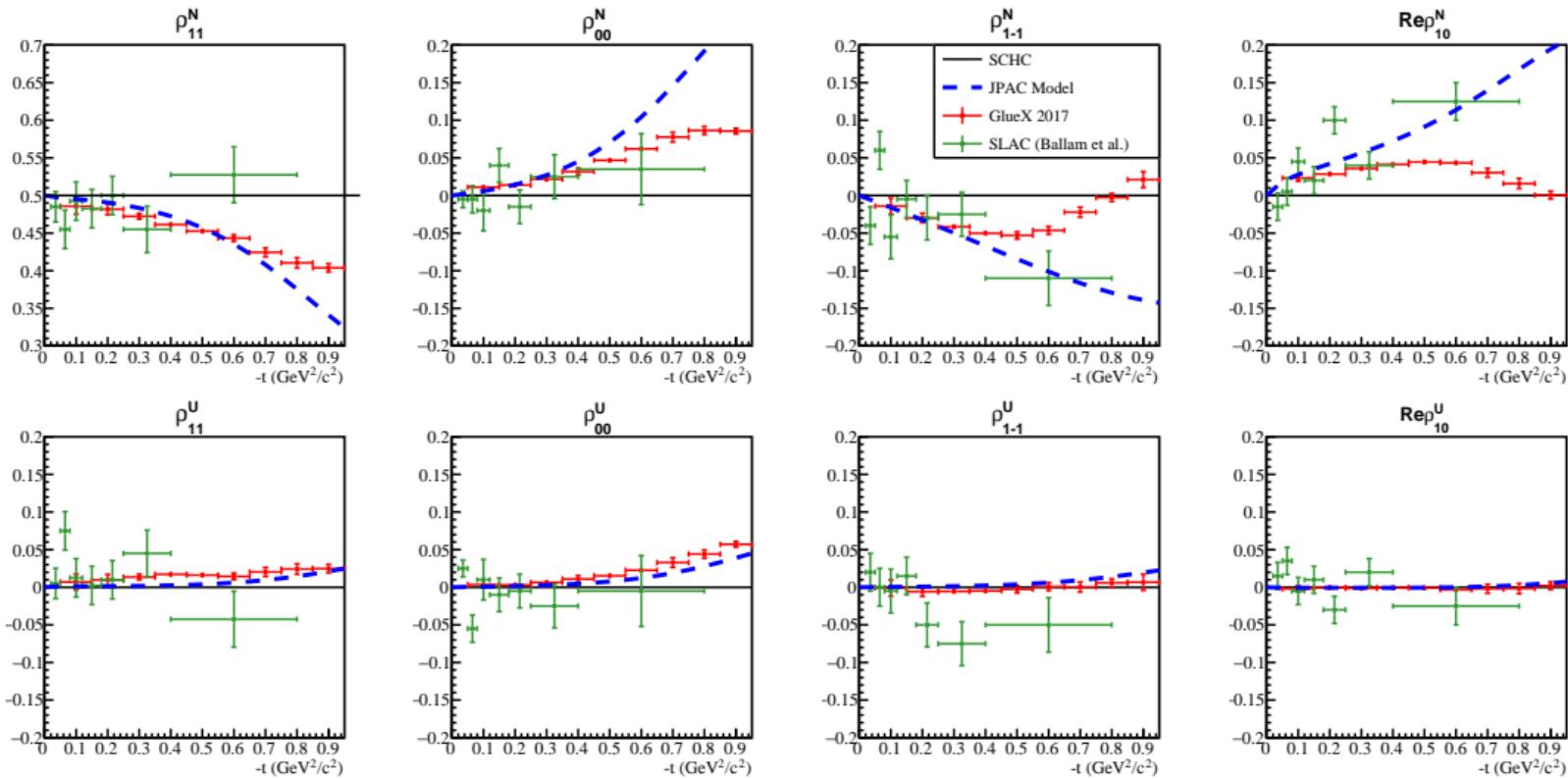
## Systematic Studies

- Analysis note revised and updated
- Presented proposal for systematic uncertainties
- Take out mass dependence from systematics, since it is physics?
- Additional sources?

## Physics Results for Paper

- Main result are SDMEs as a function of  $t$ , comparison with JPAC
- Linear combinations for natural/unnatural parity exchange contribution
- Relationship and constraints between individual SDMEs
- Mass dependence of SDMEs and interference with S-wave
- Connection to polarized reflectivity amplitudes

# (Un)Natural Parity Exchange



# SDMEs Expressed in Amplitudes



V. Mathieu, [Phys.Rev.D 100 (2019) 5, 054017]

$$\rho_{mm'}^{\alpha,\ell\ell'} = {}^{(+)}\rho_{mm'}^{\alpha,\ell\ell'} + {}^{(-)}\rho_{mm'}^{\alpha,\ell\ell'} .$$

$$\begin{aligned} {}^{(\epsilon)}\rho_{mm'}^{0,\ell\ell'} &= \kappa \sum_k \left( [\ell]_{m;k}^{(\epsilon)} [\ell']_{m';k}^{(\epsilon)*} \right. \\ &\quad \left. + (-1)^{m-m'} [\ell]_{-m;k}^{(\epsilon)} [\ell']_{-m';k}^{(\epsilon)*} \right), \quad (\text{D8a}) \end{aligned}$$

$$\begin{aligned} {}^{(\epsilon)}\rho_{mm'}^{1,\ell\ell'} &= -\epsilon \kappa \sum_k \left( (-1)^m [\ell]_{-m;k}^{(\epsilon)} [\ell']_{m';k}^{(\epsilon)*} \right. \\ &\quad \left. + (-1)^{m'} [\ell]_{m;k}^{(\epsilon)} [\ell']_{-m';k}^{(\epsilon)*} \right), \quad (\text{D8b}) \end{aligned}$$

$$\begin{aligned} {}^{(\epsilon)}\rho_{mm'}^{2,\ell\ell'} &= -i\epsilon \kappa \sum_k \left( (-1)^m [\ell]_{-m;k}^{(\epsilon)} [\ell']_{m';k}^{(\epsilon)*} \right. \\ &\quad \left. - (-1)^{m'} [\ell]_{m;k}^{(\epsilon)} [\ell']_{-m';k}^{(\epsilon)*} \right), \quad (\text{D8c}) \end{aligned}$$